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Assia Benabdallah, Patricia Gaitan, Jérôme Le Rousseau. Stability of discontinuous diffusion coefficients and initial conditions in an inverse problem for the heat equation. SIAM Journal on Control and Optimization, 2007, 46 (5), pp.1849-1881. 10.1137/050640047 . hal-00016490v3

HAL Id: hal-00016490

<https://hal.science/hal-00016490v3>

Submitted on 20 Oct 2006

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Stability of discontinuous diffusion coefficients and initial conditions in an inverse problem for the heat equation

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October 20, 2006

Abstract

We consider the heat equation with a discontinuous diffusion coefficient and give uniqueness and stability results for both the diffusion coefficient and the initial condition from a measurement of the solution on an arbitrary part of the boundary and at some arbitrary positive time. The key ingredient is the derivation of a Carleman-type estimate. The diffusion coefficient is assumed to be discontinuous across an interface with a monotonicity condition.

AMS 2000 subject classification: 35K05, 35R30.

0 Introduction

This article is devoted to the question of the identification of a diffusion coefficient, c , for a heat transfer problem in a bounded domain, with the main particularity that c is discontinuous. Such regularity can be encountered in the case of embedded materials.

Let $\Omega \subset \mathbb{R}^n$ be a bounded connected open set. The set $\overline{\Omega}$ is assumed to be a \mathcal{C}^2 submanifold with boundary in \mathbb{R}^n (see e.g. [12, Definition 1.2.1.2]). We set $\Gamma = \partial\Omega$. Let Ω_0 and Ω_1 be two non-empty open subsets of Ω such that

$$\Omega_0 \Subset \Omega, \text{ and } \Omega_1 = \Omega \setminus \overline{\Omega}_0.$$

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We denote by $S = \overline{\Omega}_0 \cap \overline{\Omega}_1$ the interface, which is assumed to be \mathcal{C}^2 . We shall use the following notation $\Omega' = \Omega_0 \cup \Omega_1$. It should be emphasized here that the position of the interface itself is not assumed to be known.

Let $T > 0$. We consider the following transmission problem for the heat equation:

$$(0.1) \quad \begin{cases} \partial_t y - \nabla \cdot (c \nabla y) = 0 & \text{in } (0, T) \times \Omega', \\ y(t, x) = h(t, x) & \text{on } (0, T) \times \Gamma, \\ \text{transmission conditions (TC1)} & \text{on } (0, T) \times S, \\ y(0, x) = y_0(x), & \text{in } \Omega, \end{cases}$$

with

$$(TC1) \quad y|_{[0, T] \times S_0} = y|_{[0, T] \times S_1}, \quad c_0 \partial_n y|_{[0, T] \times S_0} = c_1 \partial_n y|_{[0, T] \times S_1},$$

where $f|_{S_i}$ is the trace on S of $f|_{\Omega_i}$ and where

$$c = \begin{cases} c_0 & \text{in } \Omega_0, \\ c_1 & \text{in } \Omega_1, \end{cases} \quad \tilde{c} = \begin{cases} \tilde{c}_0 & \text{in } \Omega_0, \\ \tilde{c}_1 & \text{in } \Omega_1. \end{cases}$$

The boundary condition $h(t, x)$ shall be kept fixed. The diffusion coefficient c shall be kept independent of time, t . If we change the diffusion coefficient c into \tilde{c} we let \tilde{y} be the solution of (0.1) associated to \tilde{c} and \tilde{y}_0 for initial condition. In the case studied here, the interface remains unchanged when changing coefficients. Its position is however not known.

We assume that we can measure both the normal flux $\partial_n \partial_t y$ on $\gamma \subset \partial\Omega$ on the time interval (t_0, T) for some $t_0 \in (0, T)$ and Δy in Ω at time $T' \in (t_0, T)$. In the case of piecewise *constant* diffusion coefficients, i.e. $c|_{\Omega_i}$, $i = 0, 1$, is constant, our main results are (i) the injectivity of the map

$$\begin{aligned} L^\infty(\Omega) \times L^2(\Omega) &\rightarrow L^2((t_0, T) \times \gamma) \times L^2(\Omega), \\ (c, y_0) &\mapsto (\partial_n \partial_t y, \Delta y(T')), \end{aligned}$$

(uniqueness); (ii) the stability for the diffusion coefficient, c (Theorem 2.9): there exists $C > 0$ such that

$$|c - \tilde{c}|_{L^\infty(\Omega)}^2 \leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0, T) \times \gamma)}^2 + C |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\Omega')}^2;$$

and (iii) the stability for the initial condition, y_0 (Theorem 4.5): there exists $C > 0$ such that

$$|y_0 - \tilde{y}_0|_{L^2(\Omega)} \leq C / \left| \ln \left(|(y - \tilde{y})(T')|_{H^2(\Omega')} + |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0, T) \times \gamma)}^2 \right) \right|,$$

for $|(y - \tilde{y})(T')|_{H^2(\Omega')} + |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0, T) \times \gamma)}^2$ sufficiently small. At the end of Section 2, we shall observe that we may only require the difference of the coefficients, $c - \tilde{c}$, to be piecewise constant.

The key ingredient to these stability results is a *global* Carleman estimate for the operator $\partial_t - \nabla \cdot (c \nabla (\cdot))$ and the open set Ω .

The use of Carleman estimates to achieve uniqueness and stability results in inverse problems is now well-established. Some authors make use of *local* Carleman inequalities and deduce uniqueness and Hölder estimates (see [18],[17] and references cited therein). Others make use of *global* Carleman inequalities and deduce Lipschitz stability results (and hence uniqueness results). We shall follow this second approach. To our knowledge, this method was first used in [19] and then by others, e.g. [2]. For literature on Carleman estimates we refer to [13, Chapter VIII], [14, Section: 28.2–3] for local estimates and [16] for the parabolic case. For global estimates we refer to [10] and [9].

Stability results for parabolic equations are recent, to our knowledge, (see [18],[8]). Apart from [19] there are few results on Lipschitz stability, even for linear cases.

One of the main difficulties in the present problem is to deal with discontinuous diffusion coefficients. Controllability for such parabolic equations has been studied by [7]. Null-controllability property is proved via an observability inequality for the adjoint system, which is deduced in [7] from a *global* Carleman estimate yet assuming a *monotonicity* on the coefficients c in connection to the observation location: roughly speaking, the observation zone has to be located in the region where the diffusion coefficient is the smallest. Here, to achieve a stability result we have to derive a Carleman estimate for the difference of the two solutions, y, \tilde{y} . This difference is solution of a non-homogeneous parabolic equation (with discontinuous coefficient); because of the discontinuity of the diffusion coefficients it does not satisfy the appropriate transmission conditions (TC1), on the interfaces S , defined above. For this reason, under the same monotonicity assumption as in [7], we derive a peculiar Carleman estimate which includes additional *interface terms* (see Theorem 1.2).

To obtain a stability result, one has to ‘manage’ the dependence of some constants with respect to (w.r.t.) the parameters, s and λ , that appear in the weight functions used in the Carleman estimate (see (1.4) in Section 1). The interface terms require some careful treatment. In particular, a stationary-phase argument is used to obtain a sufficiently sharp asymptotic estimates of these terms for s and λ large. Usually, stability estimate are obtained by letting the parameter s become large. Here we also make use of the second parameter λ (see Section 3).

As we are concerned with parabolic equations, we have to assume the observation of the solution occurs at some positive time, $T' > 0$. The suppression of this assumption remains an open problem and it appears in all articles deriving Lipschitz stability estimates from global Carleman inequalities (see the discussion in the introduction of [19]). At the end of section 2 we show that if the position of the interface S is known, we can however localize in space the observation at time T' .

The article is organized as follows. In Section 1 we derive a Carleman estimate adapted to our problem. In Section 2 we prove a stability result for the piecewise constant diffusion coefficient c when one of the solutions, say \tilde{y} , is in a particular class of solutions. In Section 3 we prove that this class is non empty. As mentioned above we also slightly relax the piecewise-constant condition by solely imposing that the difference of the coefficients be piecewise constant. In

Section 4 we prove a stability result for the initial condition under some additional assumptions, in particular on the initial condition itself. Section 5 is devoted to a generalization of the results to a more complicated geometry, for instance allowing for more than two embedded materials. Appendix A provides some basic regularity properties for the solutions to parabolic equation with non-smooth coefficients and provides a technical lemma.

We now give some notations and important assumptions. We denote by n the outward unit normal to Ω_1 on S and also the outward unit normal to Ω on Γ . Let S_0 (resp. S_1) be the side of the interface S corresponding to the positive (resp. negative) direction of the normal n .

Note that we do not assume that Ω_0 nor Ω_1 is a connected open set. We shall however assume that they are formed with a finite number of connected open sets, say $\Omega_{0,1}, \dots, \Omega_{0,p_0}$, and $\Omega_{1,1}, \dots, \Omega_{1,p_1}$, $p_0, p_1 \in \mathbb{N}$. We shall then denote by S_{ij} the interface (possibly empty) between $\Omega_{0,i}$ and $\Omega_{1,j}$.

We make the following assumption

Assumption 0.1. *The diffusion coefficient satisfies $c_i = c|_{\Omega_i} \in \mathcal{C}^1(\overline{\Omega_i})$, $i = 0, 1$, and is independent of time t .*

Assumption 0.2. $c_0|_S \geq c_1|_S$ and $0 < c_{\min} \leq c(x) \leq c_{\max}$, $x \in \Omega'$.

Remark 0.3. Assumption 0.1 will be significantly strengthened in Section 2 to obtain a stability result. Yet, for some of the results such as the Carleman estimate proved in Section 1 and the regularity properties proved in Section 3, which can be of some use elsewhere, Assumption 0.1 is sufficient.

We let γ be a subset of the boundary Γ satisfying

Assumption 0.4. *The interior of γ is non-empty with respect to the topology on Γ induced by the Euclidean topology on \mathbb{R}^n . Each component of Ω_1 contains part of the interior of γ in its boundary.*

Examples of situations in which Assumption 0.4 is satisfied are given in Figure 1.

To obtain a Carleman estimate we introduce a geometric assumption, following [7].

Assumption 0.5. Geometric Condition (GC)

We assume that there exist two disjoint open subsets $\mathcal{O}^{(1)}, \mathcal{O}^{(2)} \Subset \Omega_0$ and two vector fields, $\zeta^{(i)} \in \mathcal{C}^1(\overline{\Omega_0}, \mathbb{R}^2)$, $i = 1, 2$, such that

$$\zeta^{(i)}(x) \cdot n(x) > 0, \quad \forall x \in S, \quad i = 1, 2,$$

$$\zeta^{(i)}(x) \cdot n(x) > 0, \quad \forall x \in \partial \mathcal{O}^{(i)}, \quad i = 1, 2,$$

$$\zeta^{(i)}(x) \neq 0, \quad \forall x \in \Omega_0 \setminus \mathcal{O}^{(i)}, \quad i = 1, 2$$

(n is the outward unit normal to Ω_1 on S and the inward unit normal to $\mathcal{O}^{(i)}$ on $\partial \mathcal{O}^{(i)}$, $i = 1, 2$). Let $x^{(i)}$ be the integral curves of $\zeta^{(i)}$, i.e.

$$\begin{cases} \frac{dx^{(i)}(t)}{dt} = \zeta^{(i)}(x^{(i)}(t)), & t > 0, \\ x^{(i)}(0) = x_0, & x_0 \in S. \end{cases}$$

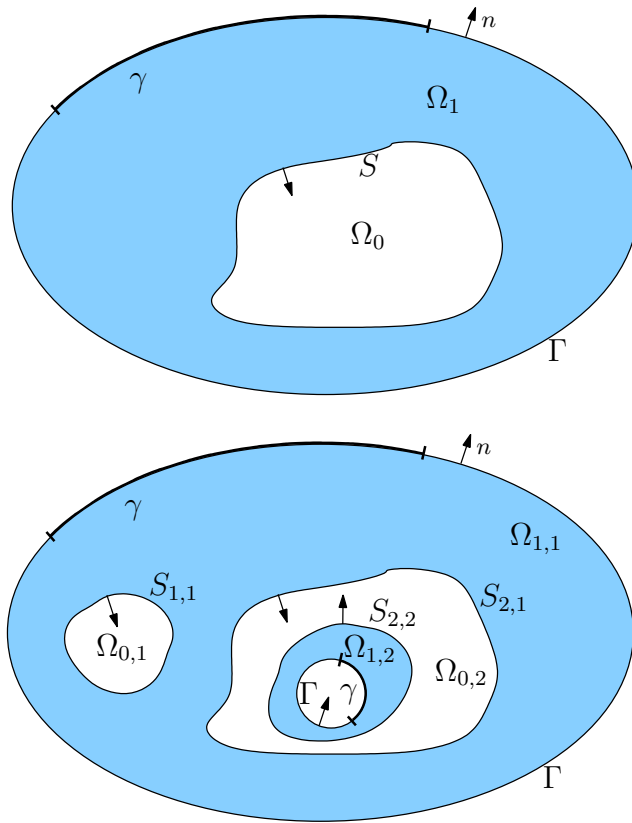


Figure 1: Geometric situations in which Assumption 0.4 and the geometric condition (GC) are satisfied. Shaded is Ω_1 . Arrows represent the normal unit vector n .

We also assume that there exists $\mathcal{T} > 0$ such that for all $x_0 \in S$, there exists $t^{(i)}(x_0) < \mathcal{T}$ satisfying

$$\begin{aligned} x^{(i)}(t) &\in \Omega_0 \setminus \mathcal{O}^{(i)}, \quad \text{for } 0 < t < t^{(i)}(x_0), \quad x_0 \in S, \quad i = 1, 2, \\ x^{(i)}(t^{(i)}(x_0)) &\in \partial \mathcal{O}^{(i)}, \quad \text{for } x_0 \in S, \quad i = 1, 2. \end{aligned}$$

Note that in Assumption 0.5, there is no restriction to having Ω_0 composed with p_0 components. The examples given in Figure 1 satisfy Assumption 0.5.

We denote by $W^{m,p}(\Omega)$, $m \in \mathbb{N}$, $1 \leq p \leq \infty$, the usual Sobolev space defined by

$$W^{m,p} = \{u \in L^p(\Omega); \partial^\alpha u \in L^p(\Omega) \text{ for } |\alpha| \leq m\},$$

where $\alpha = (\alpha_1, \dots, \alpha_n)$ is a multi-index and differentiation is to be understood in the weak sense. As usual we write $H^m(\Omega) = W^{m,2}(\Omega)$. For the definition of $W^{r,p}$, for $r \in \mathbb{R} \setminus \mathbb{N}$ we refer for instance to [1].

1 A Carleman estimate

We prove here a Carleman-type estimate with a boundary term on γ in the r.h.s. of the estimate. For this purpose we shall first introduce a particular type of weight functions, which are constructed using the following lemma.

Lemma 1.1. *Assume that there exist two disjoint open subsets $\mathcal{O}^{(1)}, \mathcal{O}^{(2)} \Subset \Omega_0$ satisfying (GC). Let γ be a subset of $\Gamma = \partial\Omega$ satisfying Assumption 0.4 and $B^{(i)}$ and $\tilde{B}^{(i)}$, $i = 1, 2$ be open balls such that $B^{(1)} \Subset \tilde{B}^{(1)} \Subset \mathcal{O}^{(1)}$ and $B^{(2)} \Subset \tilde{B}^{(2)} \Subset \mathcal{O}^{(2)}$. Then there exists two functions $\tilde{\beta}^{(1)}$ and $\tilde{\beta}^{(2)}$ such that*

$$\tilde{\beta}^{(1)}(x) = \begin{cases} \tilde{\beta}_0^{(1)} & \text{in } \Omega_0, \\ \tilde{\beta}_1 & \text{in } \bar{\Omega}_1, \end{cases} \quad \tilde{\beta}^{(2)}(x) = \begin{cases} \tilde{\beta}_0^{(2)} & \text{in } \Omega_0, \\ \tilde{\beta}_1 & \text{in } \bar{\Omega}_1, \end{cases}$$

and the functions $\tilde{\beta}_0^{(1)}, \tilde{\beta}_0^{(2)}$, and $\tilde{\beta}_1$ satisfy the following properties: $\tilde{\beta}_1 \in \mathcal{C}^2(\bar{\Omega}_1)$, $\tilde{\beta}_1 > 0$ in Ω_1 , and

$$\begin{aligned} \tilde{\beta}_1 &= 0 \text{ on } \Gamma \setminus \gamma, \quad \partial_n \tilde{\beta}_1 < 0 \text{ on } \Gamma \setminus \gamma, \\ \tilde{\beta}_1 &= 2 \text{ on } S, \quad \partial_n \tilde{\beta}_1 < 0 \text{ on } S, \end{aligned}$$

and

$$|\nabla \tilde{\beta}_1| > 0 \text{ in } \bar{\Omega}_1;$$

for $i = 1, 2$, $\tilde{\beta}_0^{(i)} \in \mathcal{C}^2(\bar{\Omega}_0)$, $\tilde{\beta}_0^{(i)} > 0$ in Ω_0 ,

$$\tilde{\beta}_0^{(i)} = \tilde{\beta}_1 = 2 \text{ on } S, \quad i = 1, 2,$$

$$c_0 \partial_n \tilde{\beta}_0^{(i)} = c_1 \partial_n \tilde{\beta}_1 \text{ on } S, \quad i = 1, 2,$$

$$(1.1) \quad \tilde{\beta}_0^{(1)} \geq 2\tilde{\beta}_0^{(2)} \text{ in } \tilde{B}^{(2)},$$

$$(1.2) \quad \tilde{\beta}_0^{(2)} \geq 2\tilde{\beta}_0^{(1)} \text{ in } \tilde{B}^{(1)},$$

and

$$(1.3) \quad |\nabla \tilde{\beta}_0^{(i)}| > 0 \text{ in } \bar{\Omega}_0 \setminus B^{(i)}, \quad i = 1, 2.$$

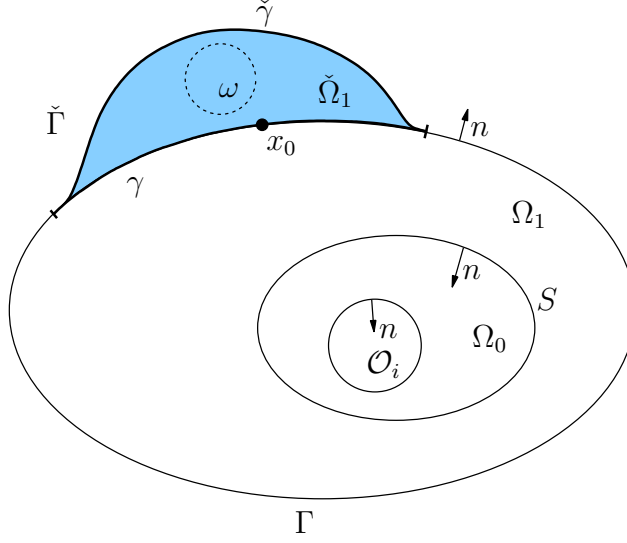


Figure 2: Geometrical situation for the proof Lemma 1.1. The shaded area represents $\check{\Omega}_1 \setminus \Omega_1 = \check{\Omega} \setminus \Omega$.

Proof. For the construction of $\tilde{\beta}_0^{(i)}$, $i = 1, 2$, supported in the connected components of Ω_0 , we refer to [7, Lemma 3.2]. We briefly show how the function $\tilde{\beta}_1$ is constructed. In actuality, the procedure described here has to be performed in each connected component of Ω_1 , which is possible since each component contains part of the interior of γ in its boundary by Assumption 0.4.

Let x_0 be in the interior of γ . We can enlarge the open set Ω_1 locally around x_0 while preserving the \mathcal{C}^2 regularity of the boundary. Such a procedure is performed in a neighborhood U of x_0 such that $U \cap \Gamma \subset \gamma$. (This can be done by locally straightening out the boundary γ as $\bar{\Omega}$ is assumed to be a \mathcal{C}^2 submanifold with boundary in \mathbb{R}^n [12, Definition 1.2.1.2]). This enlarging procedure only affects γ and leaves $\Gamma \setminus \gamma$ untouched. We call the new boundary $\check{\Gamma}$. We denote $\check{\Omega}_1$ the extension of Ω_1 and $\check{\Omega}$ that of Ω ($\Omega_1 \subset \check{\Omega}_1$, $\Omega \subset \check{\Omega}$ and $\check{\Gamma} = \partial\check{\Omega}$). Let ω be an open subset such that $\omega \Subset \check{\Omega}_1 \setminus \Omega_1$. The geometry we describe here is illustrated in Figure 2. Following [7, 10], there exists $\mu \in \mathcal{C}^2(\bar{\check{\Omega}})$ that satisfies

$$\begin{aligned} \mu &= 0, \quad \partial_n \mu < 0 \text{ on } \check{\Gamma}, \\ \mu &= 2, \quad c_0 \partial_n \tilde{\beta}_0^{(i)} = c_1 \partial_n \mu < 0 \text{ on } S, \quad i = 1, 2. \\ |\nabla \mu| &> 0 \text{ in } \check{\Omega}_1 \setminus \omega \end{aligned}$$

The function $\tilde{\beta}_1 := \mu|_{\bar{\Omega}_1}$ satisfies the required properties. ■

Choosing two functions $\tilde{\beta}^{(i)}$, $i = 1, 2$, as in the previous lemma, we introduce $\beta^{(i)} = \tilde{\beta}^{(i)} + K$ with $K = m\|\tilde{\beta}^{(1)}\|_\infty = m\|\tilde{\beta}^{(2)}\|_\infty$ and $m > 1$. For $\lambda > 0$ and

$t \in (t_0, T)$, we define the following weight functions

$$(1.4) \quad \varphi^{(i)}(t, x) = \frac{e^{\lambda\beta^{(i)}(x)}}{(t-t_0)(T-t)}, \quad \eta^{(i)}(t, x) = \frac{e^{\lambda\bar{\beta}} - e^{\lambda\beta^{(i)}(x)}}{(t-t_0)(T-t)}, i = 1, 2,$$

with $\bar{\beta} = 2m\|\tilde{\beta}^{(i)}\|_\infty$, $i = 1, 2$ (see [7],[9]). We let $t_0 \in (0, T)$ and we set $Q = (t_0, T) \times \Omega$, $Q' = (t_0, T) \times \Omega'$ and recall that $\Omega' = \Omega_0 \cup \Omega_1$.

Let $g \in H^1([t_0, T], H^{\frac{1}{2}}(S))$. We introduce transmission conditions (TC2) on the interval $[t_0, T]$ which reads

$$(TC2) \quad \begin{aligned} q|_{[t_0, T] \times S_0} &= q|_{[t_0, T] \times S_1}, \\ c_0 \partial_n q|_{[t_0, T] \times S_0} &= c_1 \partial_n q|_{[t_0, T] \times S_1} + g(t, x), \end{aligned}$$

for a function q which is H^2 in each open set Ω_i , $i = 0, 1$.

We introduce

$$\begin{aligned} \mathfrak{N}_g = \Big\{ q \in H^1(t_0, T, H_0^1(\Omega)); q|_{(t_0, T) \times \Omega_i} \in L^2(t_0, T, H^2(\Omega_i)), i = 0, 1, \\ \text{and } q \text{ satisfies (TC2) a.e. w.r.t. } t \Big\}. \end{aligned}$$

Theorem 1.2. *Let γ be a subset of the boundary Γ of an open set Ω of \mathbb{R}^n that satisfies Condition (GC), and γ satisfies Assumption 0.4. Let c satisfy Assumptions 0.1 and 0.2. Assume further that $c_0|_S - c_1|_S \geq \Delta > 0$. Let $g \in H^1(t_0, T, H^{\frac{1}{2}}(S))$. There exists $\lambda_1 = \lambda_1(\Omega, \gamma, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, c_{min}, c_{max}, \Delta) > 0$, $s_1 = s_1(\lambda_1) > 0$ and a positive constant $C = C(\Omega, \gamma, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, c_{min}, c_{max}, \Delta)$ so that the following estimate holds*

$$\begin{aligned} (1.5) \quad & |M_1^{(1)}(e^{-s\eta^{(1)}} q)|_{L^2(Q')}^2 + |M_1^{(2)}(e^{-s\eta^{(2)}} q)|_{L^2(Q')}^2 \\ & + |M_2^{(1)}(e^{-s\eta^{(1)}} q)|_{L^2(Q')}^2 + |M_2^{(2)}(e^{-s\eta^{(2)}} q)|_{L^2(Q')}^2 \\ & + s\lambda^2 \iint_Q (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |\nabla q|^2 dx dt \\ & + s^3 \lambda^4 \iint_Q (e^{-2s\eta^{(1)}} \varphi^{(1)3} + e^{-2s\eta^{(2)}} \varphi^{(2)3}) |q|^2 dx dt \\ & \leq C \left[s\lambda \int_{t_0}^T \int_\gamma (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |\partial_n q|^2 d\sigma dt \right. \\ & + \iint_{Q'} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t q \pm \nabla \cdot (c \nabla q)|^2 dx dt \\ & + s\lambda \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |g|^2 d\sigma dt \\ & + \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)4} + e^{-2s\eta^{(2)}} \varphi^{(2)4}) |g|^2 d\sigma dt \\ & \left. + s^{-2} \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t g|^2 d\sigma dt \right], \end{aligned}$$

for $s \geq s_1$, $\lambda \geq \lambda_1$ and for all $q \in \mathfrak{N}_g$, with M_1 and M_2 to be defined in (1.9)–(1.10).

We recall that $\Omega' = \Omega_0 \cup \Omega_1$ and $Q' = \Omega' \times (0, T)$. For a function ρ with a trace on the interface S , from both sides, defined in some sense, we shall denote ρ_i the trace of $\rho|_{\Omega_i}$ on S , $i = 0, 1$, when there is no ambiguity; in the case $\rho_0 = \rho_1$ we shall simply write ρ . We shall use the notation $[\rho]_S = \rho_0 - \rho_1$ for the jump of ρ across the interface S . We shall adopt Einstein's summation convention for repeated indices.

Proof. We consider $s > 0$ and $q \in \mathfrak{N}_g$. Let us set $f = \partial_t q + \nabla \cdot (c \nabla q)$ (we treat the case of the operator $\partial_t + \nabla \cdot (c \nabla)$ in the proof; the other case can be treated similarly). Then $f \in L^2(Q')$ (because of transmission conditions (TC2), observe that q is not in the domain of the operator $\nabla \cdot (c \nabla)$).

In the first part of the proof we shall write η , φ , M_1 , etc, in place of $\eta^{(i)}$, $\varphi^{(i)}$, $M_1^{(i)}$, etc, $i = 1, 2$, and treat the two cases at a time. We set $\psi = e^{-s\eta} q$. We observe that $\psi(t_0) = \psi(T) = 0$ and since q satisfies transmission conditions (TC2) (and $q(t, \cdot)|_{\Omega_i} \in H^2(\Omega_i)$ a.e. w.r.t. t), we have (a.e. w.r.t. t)

$$(1.6) \quad c_0 \partial_n \psi_{0|S}(t, \cdot) = c_1 \partial_n \psi_{1|S}(t, \cdot) + g_s(t, \cdot) \quad \text{on } S,$$

$$(1.7) \quad \nabla_\tau \psi_{0|S}(t, \cdot) = \nabla_\tau \psi_{1|S}(t, \cdot) \quad \text{on } S,$$

$$(1.8) \quad \psi_{0|S}(t, \cdot) = \psi_{1|S}(t, \cdot) \quad \text{on } S,$$

where $g_s = e^{-s\eta} g$ and ∇_τ denotes the component of the gradient that is tangential to S .

The function ψ satisfies in each Ω_i , $i = 0, 1$,

$$M_1 \psi + M_2 \psi = f_s$$

with

$$(1.9) \quad M_1 \psi = \nabla \cdot (c \nabla \psi) + s^2 \lambda^2 \varphi^2 |\nabla \beta|^2 c \psi + s(\partial_t \eta) \psi,$$

$$(1.10) \quad M_2 \psi = \partial_t \psi - 2s \lambda \varphi c \nabla \beta \cdot \nabla \psi - 2s \lambda^2 \varphi c |\nabla \beta|^2 \psi,$$

$$(1.11) \quad f_s = e^{-s\eta} f + s \lambda \varphi \nabla \cdot (c \nabla \beta) \psi - s \lambda^2 \varphi c |\nabla \beta|^2 \psi.$$

We have

$$|M_1 \psi|_{L^2(Q')}^2 + |M_2 \psi|_{L^2(Q')}^2 + 2(M_1 \psi, M_2 \psi)_{L^2(Q')} = |f_s|_{L^2(Q')}^2$$

With the same notations as in [7, Theorem 3.3], we write $(M_1 \psi, M_2 \psi)_{L^2(Q')}$ as a sum of 9 terms I_{ij} , $1 \leq i, j \leq 3$, where I_{ij} is the inner product of the i th term in the expression of $M_1 \psi$ and the j th term in the expression of $M_2 \psi$.

As compared to the proof of the Carleman estimate in [7, Theorem 3.3] we only need to adjust the computation of I_{11} , I_{12} and I_{13} to the present case. In fact the other terms do not involve transmission conditions (1.6) in their computation and thus remain unchanged from the terms obtained in [7].

The term I_{11} follows as

$$\begin{aligned} I_{11} &= \iint_{Q'} \nabla \cdot (c \nabla \psi) \partial_t \psi \, dx dt \\ &= - \iint_{Q'} c \nabla \psi \cdot \partial_t (\nabla \psi) \, dx dt + \sum_{i=0,1} (-1)^{i+1} \int_{t_0}^T \int_S c_i \partial_n \psi_i \partial_t \psi \, d\sigma dt, \end{aligned}$$

by integration by parts; the surface integral on Γ vanishes since $\partial_t \psi = 0$ there. Noting that $\nabla \psi \cdot \partial_t(\nabla \psi) = \frac{1}{2} \partial_t(|\nabla \psi|^2)$, the first term vanishes since ψ , and thus $\nabla \psi$, vanish at $t = t_0$ and $t = T$ and c is independent of t . For the remaining surface terms we use (1.6), which yields

$$I_{11} = - \int_{t_0}^T \int_S g_s \partial_t \psi \, d\sigma dt = \int_{t_0}^T \int_S \partial_t(g_s) \psi \, d\sigma dt,$$

since $g_s \in H^1(t_0, T, H^{\frac{1}{2}}(S))$.

The term I_{12} is given by

$$\begin{aligned} I_{12} &= -2s\lambda \iint_{Q'} \varphi \nabla \cdot (c \nabla \psi) c \nabla \beta \cdot \nabla \psi \, dx dt \\ &= 2s\lambda \iint_{Q'} c \nabla \psi \cdot \nabla (\varphi c \nabla \beta \cdot \nabla \psi) \, dx dt - 2s\lambda \int_{t_0}^T \int_{\Gamma} \varphi c^2 (\nabla \beta \cdot \nabla \psi) (\partial_n \psi) \, d\sigma dt \\ &\quad + 2s\lambda \sum_{i=0,1} (-1)^i \int_{t_0}^T \int_S \varphi c_i^2 (\nabla \beta_i \cdot \nabla \psi_i) (\partial_n \psi_i) \, d\sigma dt. \end{aligned}$$

This integration by parts is justified since $\psi(t, \cdot)$ is in H^2 in each Ω_i , $i = 0, 1$. Denoting by I'_{12} the remaining volume integral, we obtain

$$\begin{aligned} I'_{12} &= 2s\lambda^2 \iint_{Q'} \varphi c^2 |\nabla \psi \cdot \nabla \beta|^2 \, dx dt + 2s\lambda \iint_{Q'} \varphi c \partial_{x_i} (c \partial_{x_j} \beta) \partial_{x_i} \psi \partial_{x_j} \psi \, dx dt \\ &\quad + s\lambda \iint_{Q'} \varphi c^2 \partial_{x_j} \beta \partial_{x_j} |\nabla \psi|^2 \, dx dt. \end{aligned}$$

We further compute the last volume integral, denoted by I''_{12} . Observe that $|\nabla \psi|_{|\Omega_i}^2$ is in $W^{1,1}(\Omega_i)$ since $\psi|_{\Omega_i}(t, \cdot) \in H^2(\Omega_i)$, $i = 0, 1$. This allows to further integrate by parts, since $(c^2 \varphi \partial_{x_j} \beta)|_{\Omega_i} \in \mathcal{C}^1(\Omega_i)$, $i = 0, 1$, and yields

$$\begin{aligned} I''_{12} &= -s\lambda \iint_{Q'} \partial_{x_j} (\varphi c^2 \partial_{x_j} \beta) |\nabla \psi|^2 \, dx dt + s\lambda \int_{t_0}^T \int_{\Gamma} \varphi c^2 \partial_n \beta |\nabla \psi|^2 \, d\sigma dt \\ &\quad + s\lambda \sum_{i=0,1} (-1)^{i+1} \int_{t_0}^T \int_S \varphi c_i^2 \partial_n \beta_i |\nabla \psi_i|^2 \, d\sigma dt. \end{aligned}$$

The remaining volume integral can be further expanded into

$$I'''_{12} = -s\lambda \iint_{Q'} \varphi \partial_{x_j} (c^2 \partial_{x_j} \beta) |\nabla \psi|^2 \, dx dt - s\lambda^2 \iint_{Q'} \varphi c^2 |\nabla \beta|^2 |\nabla \psi|^2 \, dx dt.$$

Collecting the surface integrals in a term denoted by J_{12} we find

$$\begin{aligned} I_{12} &= -s\lambda^2 \iint_{Q'} \varphi c^2 |\nabla \beta|^2 |\nabla \psi|^2 \, dx dt + 2s\lambda^2 \iint_{Q'} c^2 \varphi |\nabla \psi \cdot \nabla \beta|^2 \, dx dt \\ &\quad + X_1 + J_{12}, \end{aligned}$$

where

$$X_1 = 2s\lambda \iint_{Q'} \varphi c \partial_{x_i} (c \partial_{x_j} \beta) \partial_{x_i} \psi \partial_{x_j} \psi \, dx dt \\ - s\lambda \iint_{Q'} \varphi \partial_{x_j} (c^2 \partial_{x_j} \beta) |\nabla \psi|^2 \, dx dt.$$

We now observe that since β is constant on S we have

$$(\nabla \beta \cdot \nabla \psi_i)|_S = (\partial_n \beta \partial_n \psi_i)|_S, \quad i = 0, 1.$$

Writing $|\nabla \psi|^2 = |\nabla_\tau \psi|^2 + |\partial_n \psi|^2$ we find

$$J_{12} = s\lambda \sum_{i=0,1} (-1)^i \int_{t_0}^T \int_S \varphi c_i^2 \partial_n \beta_i |\partial_n \psi_i|^2 \, d\sigma dt \\ - s\lambda \sum_{i=0,1} (-1)^i \int_{t_0}^T \int_S \varphi c_i^2 \partial_n \beta_i |\nabla_\tau \psi_i|^2 \, d\sigma dt - s\lambda \int_{t_0}^T \int_\Gamma \varphi c^2 \partial_n \beta |\partial_n \psi|^2 \, d\sigma dt,$$

where we have used that $\psi|_\Sigma$ is constant. Recall that $\nabla_\tau \psi_0 = \nabla_\tau \psi_1$, and that $c_0 \partial_n \beta_0 = c_1 \partial_n \beta_1$ on S . From transmission conditions (TC2) we have

$$|c_0 \partial_n \psi_0|^2 = |c_1 \partial_n \psi_1|^2 + |g_s|^2 + 2(c_1 (\partial_n \psi_1) g_s), \quad \text{on } S.$$

We thus obtain

$$J_{12} = s\lambda \int_{t_0}^T \int_S \varphi [\partial_n \beta]_S |c_1 \partial_n \psi_1|^2 \, d\sigma dt \\ + s\lambda \int_{t_0}^T \int_S \varphi \partial_n \beta_0 |g_s|^2 \, d\sigma dt - s\lambda \int_{t_0}^T \int_S \varphi [c]_S (c \partial_n \beta) |\nabla_\tau \psi|^2 \, d\sigma dt \\ - s\lambda \int_{t_0}^T \int_\Gamma \varphi c^2 \partial_n \beta |\partial_n \psi|^2 \, d\sigma dt + Y_1,$$

with

$$(1.12) \quad Y_1 = 2s\lambda \int_{t_0}^T \int_S \varphi c_1 \partial_n \psi_1 \partial_n \beta_0 g_s \, d\sigma dt.$$

We thus have

$$I_{12} = -s\lambda^2 \iint_{Q'} \varphi c^2 |\nabla \beta|^2 |\nabla \psi|^2 \, dx dt + 2s\lambda^2 \iint_{Q'} c^2 \varphi |\nabla \psi \cdot \nabla \beta|^2 \, dx dt \\ + s\lambda \int_{t_0}^T \int_S \varphi [\partial_n \beta]_S |c_1 \partial_n \psi_1|^2 \, d\sigma dt + s\lambda \int_{t_0}^T \int_S \varphi \partial_n \beta_0 |g_s|^2 \, d\sigma dt \\ - s\lambda \int_{t_0}^T \int_S \varphi [c]_S (c \partial_n \beta) |\nabla_\tau \psi|^2 \, d\sigma dt - s\lambda \int_{t_0}^T \int_\Gamma \varphi c^2 \partial_n \beta |\partial_n \psi|^2 \, d\sigma dt \\ + X_1 + Y_1.$$

The term I_{13} is given by

$$\begin{aligned} I_{13} &= -2s\lambda^2 \iint_{Q'} \varphi \nabla \cdot (c \nabla \psi) c |\nabla \beta|^2 \psi \, dx dt \\ &= 2s\lambda^2 \iint_{Q'} c \nabla \psi \cdot \nabla (\varphi c |\nabla \beta|^2 \psi) \, dx dt \\ &\quad + 2s\lambda^2 \sum_{i=0,1} (-1)^i \int_{t_0}^T \int_S \varphi (c_i \partial_n \psi_i) c_i |\nabla \beta_i|^2 \psi \, d\sigma dt, \end{aligned}$$

where we have used that $\psi|_{\Gamma} = 0$. Expanding the integrand in the volume integral and using (TC2) in the surface term we obtain

$$I_{13} = 2s\lambda^2 \iint_{Q'} \varphi c^2 |\nabla \beta|^2 |\nabla \psi|^2 \, dx dt + X_2 + Y_2,$$

where

$$\begin{aligned} X_2 &= 2s\lambda^2 \iint_{Q'} \varphi c \nabla \psi \cdot \nabla (c |\nabla \beta|^2) \psi \, dx dt \\ &\quad + 2s\lambda^3 \iint_{Q'} \varphi c^2 \nabla \psi \cdot \nabla \beta |\nabla \beta|^2 \psi \, dx dt \\ &\quad + 2s\lambda^2 \int_{t_0}^T \int_S \varphi (c \partial_n \beta) [\partial_n \beta]_S (c_1 \partial_n \psi_1) \psi \, d\sigma dt, \end{aligned}$$

since $\nabla_{\tau} \beta|_S = 0$ and

$$Y_2 = 2s\lambda^2 \int_{t_0}^T \int_S \varphi c_0 (\partial_n \beta_0)^2 g_s \psi \, d\sigma dt.$$

Following the proof of Theorem 3.3 in [7] we find

$$I_{21} = \frac{1}{2} s^2 \lambda^2 \iint_{Q'} \varphi^2 c |\nabla \beta|^2 \partial_t |\psi|^2 \, dx dt = -\frac{1}{2} s^2 \lambda^2 \iint_{Q'} \partial_t (\varphi^2) c |\nabla \beta|^2 |\psi|^2 \, dx dt,$$

and

$$\begin{aligned} I_{22} &= -s^3 \lambda^3 \iint_{Q'} \varphi^3 c^2 |\nabla \beta|^2 \nabla \beta \cdot \nabla (|\psi|^2) \, dx dt \\ &= 3s^3 \lambda^4 \iint_{Q'} \varphi^3 c^2 |\nabla \beta|^4 |\psi|^2 \, dx dt + s^3 \lambda^3 \int_{t_0}^T \int_S \varphi^3 |c \partial_n \beta|^2 [\partial_n \beta]_S |\psi|^2 \, d\sigma dt + X_3, \end{aligned}$$

with X_3 given by

$$X_3 = s^3 \lambda^3 \iint_{Q'} \varphi^3 \nabla \cdot (c^2 |\nabla \beta|^2 \nabla \beta) |\psi|^2 \, dx dt.$$

The terms I_{23} , I_{31} are given by

$$I_{23} = -2s^3 \lambda^4 \iint_{Q'} \varphi^3 c^2 |\nabla \beta|^4 |\psi|^2 \, dx dt,$$

$$I_{31} = \frac{1}{2}s \iint_{Q'} \partial_t \eta \partial_t (|\psi|^2) dxdt = -\frac{1}{2}s \iint_{Q'} \partial_t^2 \eta |\psi|^2 dxdt.$$

The term I_{32} is given by

$$\begin{aligned} I_{32} &= -s^2 \lambda \iint_{Q'} \varphi (\partial_t \eta) c \nabla \beta \cdot \nabla (|\psi|^2) dxdt = s^2 \lambda^2 \iint_{Q'} \varphi (\partial_t \eta) c |\nabla \beta|^2 |\psi|^2 dxdt \\ &\quad + s^2 \lambda \iint_{Q'} \varphi \nabla \cdot ((\partial_t \eta) c \nabla \beta) |\psi|^2 dxdt, \end{aligned}$$

since $\psi|_{S_0} = \psi|_{S_1}$. Finally the term I_{33} is given by

$$I_{33} = -2s^2 \lambda^2 \iint_{Q'} \varphi c (\partial_t \eta) |\nabla \beta|^2 |\psi|^2 dxdt.$$

Collecting the terms I_{ij} just computed we obtain

$$\begin{aligned} (1.13) \quad & |M_1 \psi|_{L^2(Q')}^2 + |M_2 \psi|_{L^2(Q')}^2 + 4s \lambda^2 \iint_{Q'} c^2 \varphi |\nabla \psi \cdot \nabla \beta|^2 dxdt \\ & + 2s \lambda^2 \iint_{Q'} \varphi c^2 |\nabla \beta|^2 |\nabla \psi|^2 dxdt + 2s^3 \lambda^4 \iint_{Q'} \varphi^3 c^2 |\nabla \beta|^4 |\psi|^2 dxdt \\ & + 2s \lambda \int_{t_0}^T \int_S \varphi [\partial_n \beta]_S |c_1 \partial_n \psi_1|^2 d\sigma dt - 2s \lambda \int_{t_0}^T \int_\Gamma \varphi c^2 \partial_n \beta |\partial_n \psi|^2 d\sigma dt \\ & - 2s \lambda \int_{t_0}^T \int_S \varphi [c]_S (c \partial_n \beta) |\nabla_\tau \psi|^2 d\sigma dt + 2s^3 \lambda^3 \int_{t_0}^T \int_S \varphi^3 |c \partial_n \beta|^2 [\partial_n \beta]_S |\psi|^2 d\sigma dt \\ & \quad + 2s \lambda \int_{t_0}^T \int_S \varphi \partial_n \beta_0 |g_s|^2 d\sigma dt \\ & = |f_s|_{L^2(Q')}^2 - 2(I_{11} + X_1 + Y_1 + X_2 + Y_2 + I_{21} + X_3 + I_{31} + I_{32} + I_{33}). \end{aligned}$$

We now consider the surface terms I_{11} , Y_1 , Y_2 involving the function g_s and write

$$(1.14) \quad |I_{11}| = \left| \int_{t_0}^T \int_S \partial_t (g_s) \psi d\sigma dt \right| \leq C s^{-2} \int_{t_0}^T \int_S |\partial_t g_s|^2 d\sigma dt + C s^2 \int_{t_0}^T \int_S |\psi|^2 d\sigma dt.$$

In the proof of Lemma 1.1 we are free to choose β such that $\partial_n \beta_1 / c_0 \leq -1$. Since we assume $c_0 - c_1 \geq \Delta$ we obtain

$$(1.15) \quad [\partial_n \beta]_S = \partial_n \beta_0 - \partial_n \beta_1 = \frac{\partial_n \beta_1}{c_0} (c_1 - c_0) \geq \Delta > 0.$$

The second term in (1.14) can thus be absorbed by the term

$$2s^3 \lambda^3 \int_{t_0}^T \int_S \varphi^3 |c \partial_n \beta|^2 [\partial_n \beta]_S |\psi|^2 d\sigma dt$$

in (1.13) for s sufficiently large.

The term Y_1 in (1.12) can be estimated by

$$(1.16) \quad |Y_1| = \left| 2s\lambda \int_{t_0}^T \int_S \varphi c_1 \partial_n \psi_1 \partial_n \beta_0 g_s d\sigma dt \right| \\ \leq C_\varepsilon s\lambda \int_{t_0}^T \int_S \varphi |g_s|^2 d\sigma dt + \varepsilon s\lambda \int_{t_0}^T \int_S \varphi |c_1 \partial_n \psi_1|^2 (\partial_n \beta_0)^2 d\sigma dt, \quad \varepsilon > 0.$$

For ε sufficiently small, the second surface term in (1.16) can be ‘absorbed’ by the term

$$2s\lambda \int_{t_0}^T \int_S \varphi [\partial_n \beta]_S |c_1 \partial_n \psi_1|^2 d\sigma dt$$

in (1.13) by (1.15).

The term Y_2 can be estimated by

$$|Y_2| = \left| 2s\lambda^2 \int_{t_0}^T \int_S \varphi c_0 (\partial_n \beta_0)^2 g_s \psi d\sigma dt \right| \leq Cs\lambda \int_{t_0}^T \int_S \varphi |g_s|^2 d\sigma dt \\ + Cs\lambda^3 \int_{t_0}^T \int_S \varphi c_0^2 (\partial_n \beta_0)^4 |\psi|^2 d\sigma dt.$$

Observing that $\varphi \leq CT^4 \varphi^3$, the second surface term can be ‘absorbed’ by the term

$$2s^3 \lambda^3 \int_{t_0}^T \int_S \varphi^3 |c \partial_n \beta|^2 [\partial_n \beta]_S |\psi|^2 d\sigma dt$$

in (1.13) for s sufficiently large by (1.15). The two previous estimates are the points in the proof where the hypothesis $c_0 - c_1 \geq \Delta > 0$ is needed.

Note also that

$$s^{-2} |\partial_t g_s|^2 \leq Cs^{-2} e^{-2s\eta} |\partial_t g|^2 + C(\partial_t \eta)^2 e^{-2s\eta} |g|^2 \\ \leq Cs^{-2} e^{-2s\eta} |\partial_t g|^2 + CT^2 \varphi^4 e^{-2s\eta} |g|^2,$$

where we have used that $|\partial_t \eta| \leq CT\varphi^2$ [7, equation (90)] (which makes use of the particular choices made above for K and $\bar{\beta}$, which implies that $\bar{\beta} \leq 2\beta$).

Applying the technique presented in the proof of Theorem 3.3 in [7], the previous observations yield the following Carleman estimate (we use the notation $\eta^{(i)}$)

instead of η)

$$\begin{aligned}
(1.17) \quad & |M_1^{(i)}(e^{-s\eta^{(i)}} q)|_{L^2(Q')}^2 + |M_2^{(i)}(e^{-s\eta^{(i)}} q)|_{L^2(Q')}^2 \\
& + s\lambda^2 \iint_Q e^{-2s\eta^{(i)}} \varphi^{(i)} |\nabla q|^2 dx dt + s^3 \lambda^4 \iint_Q e^{-2s\eta^{(i)}} \varphi^{(i)3} |q|^2 dx dt \\
& \leq C \left[s\lambda \int_{t_0}^T \int_{\gamma} e^{-2s\eta^{(i)}} \varphi^{(i)} |\partial_n q|^2 d\sigma dt + s^3 \lambda^4 \int_{t_0}^T \int_{\tilde{B}^{(i)}} e^{-2s\eta^{(i)}} \varphi^{(i)3} |q|^2 dx dt \right. \\
& + \iint_{Q'} e^{-2s\eta^{(i)}} |\partial_t q - \nabla \cdot (c \nabla q)|^2 dx dt + s^{-2} \int_{t_0}^T \int_S e^{-2s\eta^{(i)}} |\partial_t g|^2 d\sigma dt \\
& \left. + \int_{t_0}^T \int_S e^{-2s\eta^{(i)}} \varphi^{(i)4} |g|^2 d\sigma dt + s\lambda \int_{t_0}^T \int_S e^{-2s\eta^{(i)}} \varphi^{(i)} |g|^2 d\sigma dt \right],
\end{aligned}$$

for $i = 1, 2$, and for $\lambda \geq \lambda_0(\Omega, \gamma, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, c_{min}, c_{max}, \Delta)$ and $s \geq s_0(\lambda_0)$ (the sets $\tilde{B}^{(i)}$ were introduced in Lemma 1.1). Note that the condition $[c]_s \geq 0$ is needed to obtain the previous estimate.

Adding (1.17) for $i = 1, 2$, we deduce (1.5) with the same argumentation as in the proof of Theorem 3.4 in [7]. The terms integrated over $(t_0, T) \times \tilde{B}^{(i)}$ are absorbed by other terms using properties (1.1)–(1.3) of $\tilde{\beta}^{(i)}$, $i = 1, 2$. ■

Remark 1.3. The Carleman estimate that was just derived is peculiar because of the presence of terms integrated on the interface S . In particular, two terms involve the function g with different powers for the parameters s and λ and for the weight functions $\varphi^{(i)}$, $i = 1, 2$. This Carleman estimate is the key ingredient in the subsequent analysis. The interface terms will require some special treatment. The two parameters s and λ will also have an important role to play in the next section.

Remark 1.4. In the case $g = 0$, the previous Carleman estimate simplifies. By inspection of the proof of Theorem 1.2, observe that in the case $g = 0$, the condition $c_{0|_S} - c_{1|_S} \geq 0$ is sufficient to obtain the Carleman estimate [7]. The case of $c_{0|_S} - c_{1|_S} < 0$ remains open in the case of a dimension greater than equal to 2 (in the one-dimensional case a Carleman estimate for the heat operator, $\partial_t \pm \partial_x(c\partial_x)$, in arbitrary situations, can be found in [3, 4]).

2 Uniqueness and stability estimate for the diffusion coefficients

In this section we establish a uniqueness result for the discontinuous diffusion coefficient c as well as a stability inequality. This inequality estimates the discrepancy in the coefficients c and \tilde{c} of two materials (with the same geometry) with an upper bound given by some Sobolev norms of the difference between

the solutions y and \tilde{y} to

$$(2.1) \quad \begin{cases} \partial_t \tilde{y} - \nabla \cdot (\tilde{c} \nabla \tilde{y}) = 0, & \text{in } (0, T) \times \Omega, \\ \tilde{y}(t, x) = h(t, x), & \text{on } (0, T) \times \Gamma, \\ \text{transmission conditions } (\widetilde{\text{TC1}}), & \text{on } [0, T] \times S, \\ \tilde{y}(0) = \tilde{y}_0, \end{cases}$$

with

$$(\widetilde{\text{TC1}}) \quad \tilde{y}|_{[0, T] \times S_0} = \tilde{y}|_{[0, T] \times S_1}, \quad \tilde{c}_0 \partial_n \tilde{y}|_{[0, T] \times S_0} = \tilde{c}_1 \partial_n \tilde{y}|_{[0, T] \times S_1},$$

and

$$(2.2) \quad \begin{cases} \partial_t y - \nabla \cdot (c \nabla y) = 0, & \text{in } (0, T) \times \Omega, \\ y(t, x) = h(t, x), & \text{on } (0, T) \times \Gamma, \\ \text{transmission conditions } (\text{TC1}), & \text{on } [0, T] \times S, \\ y(0) = y_0. \end{cases}$$

The Carleman estimate proved in the previous section will be the key ingredient in the proof of such a stability estimate.

We introduce

$$\xi = c - \tilde{c} = \begin{cases} \xi_0 = c_0 - \tilde{c}_0 & \text{in } \Omega_0, \\ \xi_1 = c_1 - \tilde{c}_1 & \text{in } \Omega_1. \end{cases}$$

We set $u = y - \tilde{y}$ and $v = \partial_t u$. Then v is solution to the following problem

$$(2.3) \quad \begin{cases} \partial_t v - \nabla \cdot (c \nabla v) = \nabla \cdot (\xi \nabla \partial_t \tilde{y}), & \text{in } (0, T) \times \Omega', \\ v = 0, & \text{on } (0, T) \times \Gamma, \\ \text{transmission conditions } (\text{TC2}), & \text{on } [0, T] \times S, \end{cases}$$

with

$$(\text{TC2}) \quad \begin{cases} v|_{[0, T] \times S_0} = v|_{[0, T] \times S_1}, \\ c_0 \partial_n v|_{[0, T] \times S_0} = c_1 \partial_n v|_{[0, T] \times S_1} + g(t, x), \end{cases}$$

where

$$g(t, x) = \xi_1 \partial_n \partial_t \tilde{y}|_{[0, T] \times S_1} - \xi_0 \partial_n \partial_t \tilde{y}|_{[0, T] \times S_0} = \alpha \partial_n \partial_t \tilde{y}|_{[0, T] \times S_0},$$

$$\text{with } \alpha = \xi_1|_{S_1} \frac{\tilde{c}_0|_{S_0}}{\tilde{c}_1|_{S_1}} - \xi_0|_{S_0}.$$

Let $T' = \frac{1}{2}(T + t_0)$. We make the following assumption.

Assumption 2.1. *The solutions \tilde{y} and y belong to $H^2(t_0, T, H^1(\Omega))$ and are such that $y|_{\Omega_i} \in H^1(t_0, T, H^2(\Omega_i))$, $\tilde{y}|_{\Omega_i} \in H^2(t_0, T, H^2(\Omega_i))$, $i = 0, 1$. Furthermore, \tilde{y} satisfies*

1. *Let $r > 0$. The solution \tilde{y} is such that $|\Delta \tilde{y}(T')| \geq r > 0$ in Ω' ;*

2. $\tilde{y}|_{\Omega_i}$ is in a bounded domain of $W^{2,\infty}(t_0, T, H^2(\Omega_i))$, $i = 0, 1$: there exists $M > 0$ such that

$$|\tilde{y}|_{\Omega_i}(t, \cdot)|_{H^2(\Omega_i)}^2 + |\partial_t \tilde{y}|_{\Omega_i}(t, \cdot)|_{H^2(\Omega_i)}^2 + |\partial_t^2 \tilde{y}|_{\Omega_i}(t, \cdot)|_{H^2(\Omega_i)}^2 \leq M, \quad i = 0, 1,$$

a.e. for $t \in (t_0, T)$;

3. $\Delta \partial_t \tilde{y}|_{\Omega_i}$ is in a bounded domain of $L^2(t_0, T, L^\infty(\Omega_i))$, $i = 0, 1$: there exists $K > 0$ such that

$$\int_{t_0}^T |\Delta \partial_t \tilde{y}|_{\Omega_i}(t, \cdot)|_{L^\infty(\Omega_i)}^2 dt \leq K^2, \quad i = 0, 1.$$

In Section 3 we shall show that for *any* initial conditions y_0, \tilde{y}_0 in $L^2(\Omega)$ we can achieve the properties listed in Assumption 2.1 by using some particular boundary conditions $h(t, x)$.

From Assumption 2.1, the functions \tilde{y} and v are such that $\tilde{y}|_{\Omega_i}, v|_{\Omega_i} \in H^2(\Omega_i)$, $i = 0, 1$. Then $g \in H^1(t_0, T, H^{\frac{1}{2}}(S))$. The second equality in condition transmission (TC2) thus takes place in the space $H^{\frac{1}{2}}(S)$. Observe that $v = \partial_t(y - \tilde{y}) \in \mathfrak{N}_g$ from the above assumption. We can thus apply Carleman estimate (1.5) to v .

We shall use the notations of the proof of Theorem 1.2. We set $\psi^{(i)} = e^{-s\eta^{(i)}} v$, $i = 1, 2$. With the operator $M_2^{(i)}$ defined in (1.10) we introduce, following [2],

$$I^{(i)} = \int_{t_0}^{T'} \int_{\Omega'} M_2^{(i)} \psi^{(i)} \varphi^{(i)\frac{3}{2}} \psi^{(i)} dx dt, \quad i = 1, 2, \quad \text{and } I = \frac{1}{2}(I^{(1)} + I^{(2)}).$$

Note the additional $\varphi^{(i)\frac{3}{2}}$ factor as compared to [2]. This will be of importance below.

We have the following estimates.

Lemma 2.2. *Let $\lambda \geq \lambda_1$ and $s \geq s_1$ then*

$$\begin{aligned} |I| \leq & C s^{-3/2} \lambda^{-2} \left[s \lambda \int_{t_0}^T \int_{\gamma} (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |\partial_n v|^2 d\sigma dt \right. \\ & + \int \int_{Q'} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t v - \nabla \cdot (c \nabla v)|^2 dx dt \\ & + s \lambda \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |g|^2 d\sigma dt \\ & + \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)4} + e^{-2s\eta^{(2)}} \varphi^{(2)4}) |g|^2 d\sigma dt \\ & \left. + s^{-2} \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t g|^2 d\sigma dt \right]. \end{aligned}$$

Proof. Observe that

$$|I^{(i)}| \leq \frac{1}{2} s^{-3/2} \lambda^{-2} \left(|M_2^{(i)} \psi^{(i)}|_{L^2(Q')}^2 + s^3 \lambda^4 \int \int_Q \varphi^{(i)3} e^{-2s\eta^{(i)}} |v|^2 dx dt \right),$$

$i = 1, 2.$

(We recall that $Q' = (t_0, T) \times \Omega'$.) Thus

$$|I| \leq \frac{1}{2} s^{-3/2} \lambda^{-2} \left(|M_2^{(1)} \psi^{(1)}|_{L^2(Q')}^2 + |M_2^{(2)} \psi^{(2)}|_{L^2(Q')}^2 \right. \\ \left. + s^3 \lambda^4 \iint_Q (e^{-2s\eta^{(1)}} \varphi^{(1)3} + e^{-2s\eta^{(2)}} \varphi^{(2)3}) |v|^2 dx dt \right),$$

which yields the result from Carleman estimate (1.5). \blacksquare

Lemma 2.3. *Let $\lambda \geq \lambda_1$ and $s \geq s_1$ then*

$$\int_{\Omega'} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}} \varphi^{(2)\frac{3}{2}} \right) (T', x) |v(T', \cdot)|^2 dx \\ \leq C s^{-3/2} \lambda^{-2} \left[s\lambda \int_{t_0}^T \int_{\gamma} (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |\partial_n v|^2 d\sigma dt \right. \\ + \iint_{Q'} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\nabla \cdot (\xi \nabla \partial_t \tilde{y})|^2 dx dt \\ + s\lambda \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |g|^2 d\sigma dt \\ + \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)4} + e^{-2s\eta^{(2)}} \varphi^{(2)4}) |g|^2 d\sigma dt \\ \left. + s^{-2} \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t g|^2 d\sigma dt \right].$$

Proof. We evaluate integral $I^{(i)}$, $i = 1, 2$, using (1.10)

$$I^{(i)} = \int_{t_0}^{T'} \int_{\Omega'} \left(\partial_t \psi^{(i)} - 2s\lambda \varphi^{(i)} c \nabla \beta^{(i)} \cdot \nabla \psi^{(i)} \right. \\ \left. - 2s\lambda^2 \varphi^{(i)} c |\nabla \beta^{(i)}|^2 \psi^{(i)} \right) \varphi^{(i)\frac{3}{2}} \psi^{(i)} dx dt \\ = \frac{1}{2} \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{3}{2}} \partial_t |\psi^{(i)}|^2 dx dt - s\lambda \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{5}{2}} c \nabla \beta^{(i)} \cdot \nabla |\psi^{(i)}|^2 dx dt \\ - 2s\lambda^2 \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{5}{2}} c |\nabla \beta^{(i)}|^2 |\psi^{(i)}|^2 dx dt \\ = \frac{1}{2} \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{3}{2}} \partial_t |\psi^{(i)}|^2 dx dt + s\lambda \int_{t_0}^{T'} \int_{\Omega'} \nabla \cdot (\varphi^{(i)\frac{5}{2}} c \nabla \beta^{(i)}) |\psi^{(i)}|^2 dx dt \\ - 2s\lambda^2 \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{5}{2}} c |\nabla \beta^{(i)}|^2 |\psi^{(i)}|^2 dx dt,$$

by integration by parts, without any remaining integral over $(t_0, T') \times S$ by condition transmission (1.8). With an integration by parts w.r.t. t in the first

integral, we then obtain

$$\begin{aligned}
(2.4) \quad & \frac{1}{2} \int_{\Omega'} \varphi^{(i)\frac{3}{2}} |\psi^{(i)}(T', \cdot)|^2 dx = I^{(i)} + \frac{1}{2} s \lambda^2 \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{5}{2}} c |\nabla \beta^{(i)}|^2 |\psi^{(i)}|^2 dx dt \\
& - s \lambda \int_{t_0}^{T'} \int_{\Omega'} \varphi^{(i)\frac{5}{2}} \nabla \cdot (c \nabla \beta^{(i)}) |\psi^{(i)}|^2 dx dt + \frac{3}{4} \int_{t_0}^{T'} \int_{\Omega'} (\partial_t \varphi^{(i)}) \varphi^{(i)\frac{1}{2}} |\psi^{(i)}|^2 dx dt, \\
& i = 1, 2,
\end{aligned}$$

since $\varphi^{(i)\frac{3}{2}} \psi^{(i)}(t_0, \cdot) = 0$. Adding (2.4) for $i = 1, 2$ we obtain

$$\begin{aligned}
(2.5) \quad & \int_{\Omega'} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}} \varphi^{(2)\frac{3}{2}} \right) (T', x) |v(T', \cdot)|^2 dx \leq 4|I| \\
& + C(s\lambda^2 + s\lambda + 1) \int_{t_0}^{T'} \int_{\Omega'} \left(e^{-2s\eta^{(1)}(t,x)} \varphi^{(1)\frac{5}{2}} + e^{-2s\eta^{(2)}(t,x)} \varphi^{(2)\frac{5}{2}} \right) |v|^2 dx dt,
\end{aligned}$$

observing that $|\partial_t \varphi^{(i)}| \leq CT \varphi^{(i)2}$, $i = 1, 2$. We use Carleman estimate (1.5) to obtain an upper-bound for the last term in (2.5) which yields the result by Lemma 2.2. \blacksquare

We shall now assume:

Assumption 2.4. *The diffusion coefficients c and \tilde{c} are piecewise constant, in the sense that $c|_{\Omega_i}$, resp. $\tilde{c}|_{\Omega_i}$, are constant in each connected component of Ω_i , $i = 0, 1$. We define*

$$\begin{aligned}
c_{0,j} &= c|_{\Omega_{0,j}}, \quad j = 1, \dots, p_0, \\
c_{1,j} &= c|_{\Omega_{1,j}}, \quad j = 1, \dots, p_1,
\end{aligned}$$

with similar notations for \tilde{c} and ξ .

In this case observe that, in Ω' ,

$$\begin{aligned}
v(T', x) &= c \Delta u(T', x) + \xi \Delta \tilde{y}(T', x) \\
&= \sum_{j=1}^{p_0} c_{0,j} \Delta u(T', x) \chi_{\Omega_{0,j}} + \sum_{j=1}^{p_1} c_{1,j} \Delta u(T', x) \chi_{\Omega_{1,j}} \\
&\quad + \sum_{j=1}^{p_0} \xi_{0,j} \Delta \tilde{y}(T', x) \chi_{\Omega_{0,j}} + \sum_{j=1}^{p_1} \xi_{1,j} \Delta \tilde{y}(T', x) \chi_{\Omega_{1,j}},
\end{aligned}$$

from the equation satisfied by u expressed at time T' and the definition of v above.

From Lemma 2.3, we obtain

$$\begin{aligned}
(2.6) \quad & \int_{\Omega'} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}} \varphi^{(2)\frac{3}{2}} \right) (T', x) |\xi \Delta \tilde{y}(T', x)|^2 dx \\
& \leq C \int_{\Omega'} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}} \varphi^{(2)\frac{3}{2}} \right) (T', x) |c \Delta u(T', x)|^2 dx \\
& + Cs^{-3/2} \lambda^{-2} \left[s\lambda \int_{t_0}^T \int_{\gamma} (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |\partial_n v|^2 d\sigma dt \right. \\
& \quad + \iint_{Q'} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\nabla \cdot (\xi \nabla \partial_t \tilde{y})|^2 dx dt \\
& \quad + s\lambda \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |g|^2 d\sigma dt \\
& \quad + \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)4} + e^{-2s\eta^{(2)}} \varphi^{(2)4}) |g|^2 d\sigma dt \\
& \quad \left. + s^{-2} \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t g|^2 d\sigma dt \right].
\end{aligned}$$

From Assumption 2.1 we find that

$$|\Delta \partial_t \tilde{y}(t, x)|^2 \leq k^2(t) |\Delta \tilde{y}(T', x)|^2, \text{ in each } (t_0, T) \times \Omega_i, \quad i = 0, 1,$$

for

$$k(t) = \frac{1}{r} \sup_{i=0,1} |\Delta \partial_t \tilde{y}|_{\Omega_i}(t, \cdot)_{L^\infty(\Omega_i)}.$$

From Assumption 2.1, $k \in L^2(t_0, T)$ and $|k|_{L^2(t_0, T)} \leq K' = \frac{1}{r} K$. These observations yield the following estimation of the third term in the r.h.s. of (2.6):

$$\begin{aligned}
(2.7) \quad & \int_{t_0}^T \int_{\Omega_i} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\xi \Delta \partial_t \tilde{y}|^2 dx dt \\
& \leq K'^2 \sum_{i=0,1} \sum_{j=1}^{p_i} |\xi_{i,j}|^2 \int_{\Omega_{i,j}} \left(e^{-2s\eta^{(1)}(T', x)} + e^{-2s\eta^{(2)}(T', x)} \right) |\Delta \tilde{y}(T', x)|^2 dx,
\end{aligned}$$

where we have used that

$$e^{-2s\eta^{(i)}(t, x)} \leq e^{-2s\eta^{(i)}(T', x)}, \quad x \in \Omega, \quad t \in (t_0, T), \quad i = 1, 2.$$

Observing that $0 < C \leq \varphi^{(i)}$, $i = 1, 2$ since $\beta^{(i)} \geq 0$ and $\frac{1}{(T-t)(t-t_0)} \geq C > 0$, we obtain

$$\begin{aligned}
& \int_{t_0}^T \int_{\Omega_i} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\xi \Delta \partial_t \tilde{y}|^2 dx dt \\
& \leq K'^2 \sum_{i=0,1} \sum_{j=1}^{p_i} |\xi_{i,j}|^2 \int_{\Omega_{i,j}} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} \right. \\
& \quad \left. + e^{-2s\eta^{(2)}(T', x)} \varphi^{(2)\frac{3}{2}} \right) (T', x) |\Delta \tilde{y}(T', x)|^2 dx.
\end{aligned}$$

We now treat the interface terms that appear in the r.h.s. of the Carleman estimate. Recall that

$$g(t, x) = \alpha \partial_n \partial_t \tilde{y}|_{[0, T] \times S_0},$$

$$\alpha = \xi_1 \frac{\tilde{c}_0}{\tilde{c}_1} - \xi_0.$$

Note that $\eta(t, \cdot)$ is constant on S . We denote this constant by $\eta(t, S)$. More generally we shall denote $\rho(S)$ the value on S of a function ρ which is constant on S . We obtain

$$\begin{aligned} & \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\partial_t g|^2 d\sigma dt \\ & \leq \int_{t_0}^T (e^{-2s\eta^{(1)}(t, S)} + e^{-2s\eta^{(2)}(t, S)}) \int_S |\partial_t g|^2 d\sigma dt \\ & \leq M' |\xi|^2 \int_{t_0}^T (e^{-2s\eta^{(1)}(t, S)} + e^{-2s\eta^{(2)}(t, S)}) dt, \end{aligned}$$

from trace inequalities and from Assumption 2.1, for $M' = C_{Tr}(1 + \frac{c_{max}}{c_{min}})^2 M$, where $|\xi| = \sqrt{\xi_0^2 + \xi_1^2}$, since $|\alpha| \leq (1 + \frac{c_{max}}{c_{min}}) |\xi|$ from Assumption 0.2. The constant C_{Tr} is the constant found in the trace estimates

$$\int_S |\partial_n \tilde{\rho}|^2 \leq C_{Tr} |\rho|_{H^2(\Omega_i)}^2, \quad i = 0, 1,$$

if $\rho|_{\Omega_i} \in H^2(\Omega_i)$, $i = 0, 1$. Similarly, since $\varphi^{(i)}$, $i = 1, 2$, are constant on S , we have

$$\begin{aligned} & \int_{t_0}^T \int_S (e^{-2s\eta^{(1)}} \varphi^{(1)j} + e^{-2s\eta^{(2)}} \varphi^{(2)j}) |g|^2 d\sigma dt \\ & \leq M' |\xi|^2 \int_{t_0}^T (e^{-2s\eta^{(1)}} \varphi^{(1)j} + e^{-2s\eta^{(2)}} \varphi^{(2)j})(t, S) dt, \quad j \in \mathbb{N}. \end{aligned}$$

With

$$w_k(s, \lambda) := \int_{t_0}^T (e^{-2s\eta^{(1)}} \varphi^{(1)k} + e^{-2s\eta^{(2)}} \varphi^{(2)k})(t, S) dt, \quad k \in \mathbb{N},$$

and

$$W_{i,j}(s, \lambda) := \int_{\Omega_{i,j}} (e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}} \varphi^{(2)\frac{3}{2}})(T', x) |\Delta \tilde{y}(T', x)|^2 dx,$$

$$i = 0, 1, j = 1, \dots, p_i,$$

we thus obtain, for $\lambda \geq \lambda_1$ and $s \geq s_1$,

$$\begin{aligned}
(2.8) \quad & \sum_{i=0,1} \sum_{j=1}^{p_i} |\xi_{i,j}|^2 \left\{ (1 - CK'^2 s^{-\frac{3}{2}} \lambda^{-2}) W_{i,j}(s, \lambda) \right. \\
& \left. - CM' [s^{-\frac{1}{2}} \lambda^{-1} w_1(s, \lambda) + s^{-\frac{3}{2}} \lambda^{-2} w_4(s, \lambda) + s^{-\frac{7}{2}} \lambda^{-2} w_0(s, \lambda)] \right\} \\
& \leq C \int_{\Omega'} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}} \varphi^{(2)\frac{3}{2}} \right) (T', x) |c\Delta u(T', \cdot)|^2 dx \\
& \quad + Cs^{-\frac{1}{2}} \lambda^{-1} \int_{t_0}^T \int_{\gamma} (e^{-2s\eta^{(1)}} \varphi^{(1)} + e^{-2s\eta^{(2)}} \varphi^{(2)}) |\partial_n v|^2 d\sigma dt.
\end{aligned}$$

To obtain a stability result we need to prove that the coefficients for $|\xi_{i,j}|^2$, $i = 0, 1$, $j = 1, \dots, p_i$, can be made positive. To do so we need to understand the behavior of the integrals $w_k(s, \lambda)$ and $W_{i,j}(s, \lambda)$ as s and λ become large.

We first establish the asymptotic behavior of $w_k(s, \lambda)$. We set

$$w_k^{(i)}(s, \lambda) := \int_{t_0}^T e^{-2s\eta^{(i)}(t,S)} \varphi^{(i)k}(t, S) dt, \quad k \in \mathbb{N}.$$

Lemma 2.5. *The following estimates holds*

$$\begin{aligned}
(2.9) \quad w_k^{(i)}(s, \lambda) = e^{-2s\eta^{(i)}(T',S)} \varphi^{(i)k}(T', S) & \left\{ \frac{\sqrt{\pi} s^{-\frac{1}{2}}}{\sqrt{\phi''(T') \sqrt{e^{\lambda\bar{\beta}} - e^{\lambda\beta^{(i)}(S)}}}} \right. \\
& \left. + \mathcal{O}\left(\frac{s^{-\frac{3}{2}}}{(e^{\lambda\bar{\beta}} - e^{\lambda\beta^{(i)}(S)})^{\frac{3}{2}}}\right) \right\},
\end{aligned}$$

with $\phi(t) = \frac{1}{(T-t)(t-t_0)}$, for $s > s_2 > 0$ and $\lambda > \lambda_2 > 0$.

Proof. Let $T^{(2)}$ and $T^{(3)}$ be such that $t_0 < T^{(2)} < T' < T^{(3)} < T$. We choose $\chi_1 \in \mathcal{C}_c^\infty([t_0, T^{(2)}])$, $\chi_2 \in \mathcal{C}_c^\infty((t_0, T))$, and $\chi_3 \in \mathcal{C}_c^\infty((T^{(3)}, T])$, all three non-negative, such that $\chi_1 + \chi_2 + \chi_3 = 1$ and $\chi_1 = 1$ in a neighborhood of t_0 , $\chi_2 = 1$ in a neighborhood of T' , and $\chi_3 = 1$ in a neighborhood of T . With this partition of unity we break $w_k^{(i)}$ into three pieces: $w_k^{(i)} = w_k^{(i,1)} + w_k^{(i,2)} + w_k^{(i,3)}$ with

$$w_k^{(i,j)}(s, \lambda) := \int_{t_0}^T e^{-2s\eta^{(i)}(t,S)} \varphi^{(i)k}(t, S) \chi_j(t) dt, \quad j = 1, 2, 3.$$

The first and the third term are treated similarly. Let $s_2 > 0$ and $\lambda_2 > 0$. We set $\tau(s, \lambda, S) = s(e^{\lambda\bar{\beta}} - e^{\lambda\beta^{(i)}(S)})$. We observe

$$\begin{aligned}
w_k^{(i,1)}(s, \lambda) &= e^{k\lambda\beta^{(i)}(S)} \int_{t_0}^T e^{-2s\eta^{(i)}(t,S)} \phi^k(t) \chi_1(t) dt \\
&\leq e^{k\lambda\beta^{(i)}(S)} e^{-2(s-s_2)\eta^{(i)}(T^{(2)},S)} \int_{t_0}^T e^{-2s_2\eta^{(i)}(t,S)} \phi^k(t) \chi_1(t) dt \\
&\leq C(s_2, \lambda_2) e^{k\lambda\beta^{(i)}(S)} e^{-2(s-s_2)\eta^{(i)}(T^{(2)},S)} \\
&\leq C e^{k\lambda\beta^{(i)}(S)} e^{-2s\eta^{(i)}(T',S)} e^{-2s(\eta^{(i)}(T^{(2)},S) - \eta^{(i)}(T',S))} \\
&= e^{k\lambda\beta^{(i)}(S)} e^{-2s\eta^{(i)}(T',S)} \mathcal{O}(\tau(s, \lambda, S)^{-l}),
\end{aligned}$$

for all $l \in \mathbb{N}$, if $s > s_2$ and $\lambda > \lambda_2$, since

$$s(\eta^{(i)}(T^{(2)}, S) - \eta^{(i)}(T', S)) = \tau(s, \lambda, S)(\phi(T^{(2)}) - \phi(T')),$$

and $\phi(T^{(2)}) - \phi(T') > 0$.

For the second term $w_k^{(i,2)}$ we write

$$\begin{aligned} w_k^{(i,2)}(s, \lambda) &:= e^{k\lambda\beta^{(i)}(S)} \int_{t_0}^T e^{-2\tau(s, \lambda, S)\phi(t)} \phi^k(t) \chi_2(t) dt \\ &= e^{k\lambda\beta^{(i)}(S)} e^{-2\tau(s, \lambda, S)\phi(T')} \int_{t_0}^T e^{-2\tau(s, \lambda, S)(\phi(t) - \phi(T'))} \phi^k(t) \chi_2(t) dt. \end{aligned}$$

We then apply the following stationary phase formula [15, Theorem 7.7.5], in one dimension,

$$\begin{aligned} \left| \int u(y) e^{i\omega f(y)} dy - e^{i\omega f(y_0)} \left(\frac{\omega f''(y_0)}{2\pi i} \right)^{-\frac{1}{2}} \sum_{l < l_1} \omega^{-l} L_l u \right| \\ \leq C \omega^{-l_1 - \frac{1}{2}} \sum_{|\alpha| \leq 2l_1} \sup |\partial^\alpha u|, \quad \omega > 0, \quad u \in \mathcal{C}_c^\infty, \end{aligned}$$

where L_l is a differential operator of order $2l$ evaluated at T' , $L_0 u = u(T')$. (To obtain the additional factor $-\frac{1}{2}$ in the r.h.s. as compared to the formula given in [15, Theorem 7.7.5], simply write the formula to the order $l_1 + 1$.) This formula is valid if $\text{Im}(f) \geq 0$, $\text{Im}(f(y_0)) = 0$, $f'(y_0) = 0$, $f''(y_0) \neq 0$, and $f'(y) \neq 0$ in $\text{supp}(u) \setminus \{y_0\}$. Here the phase function $f(t) = i(\phi(t) - \phi(T'))$ is imaginary (note that $\phi''(T') > 0$), $\omega = 2\tau(s, \lambda, S)$, and $u = \phi^k \chi_2$. The stationary point is $t = T'$. With $l_1 = 1$, we obtain

$$\begin{aligned} \left| \int_{t_0}^T e^{-2\tau(s, \lambda, S)(\phi(t) - \phi(T'))} \phi^k(t) \chi_2(t) dt - \phi^k(T') \sqrt{\frac{\pi}{\phi''(T')}} \tau(s, \lambda, S)^{-\frac{1}{2}} \right| \\ \leq C \tau(s, \lambda, S)^{-3/2}. \end{aligned}$$

This yields (2.9). ■

To achieve our goal we also need an estimation from below for the terms $W_{i,j}(s, \lambda)$. We set

$$\begin{aligned} W_{k,j}^{(i)}(s, \lambda) &:= \int_{\Omega_{k,j}} e^{-2s\eta^{(i)}(T', x)} \varphi^{(i)\frac{3}{2}}(T', x) |\Delta \tilde{y}(T', x)|^2 dx, \\ k &= 0, 1, \quad j = 1, \dots, p_k, \quad i = 1, 2. \end{aligned}$$

For the terms $W_{0,j}^{(i)}(s, \lambda)$, we have the following

Lemma 2.6. *Let $\varepsilon > 0$. We have*

$$\begin{aligned} W_{0,j}^{(i)}(s, \lambda) &\geq C_{s_2, i, j} \frac{r^2 |S_j|}{s\lambda} e^{-2s\eta^{(i)}(T', S)} (\varphi^{(i)}(T', S))^{\frac{1}{2}} e^{-\lambda\varepsilon}, \\ i &= 1, 2, \quad j = 1, \dots, p_0, \end{aligned}$$

for $s \geq s_2 > 0$, $\lambda > 0$ and where $S_j = \bigcup_{k=1, \dots, p_1} S_{jk}$.

Proof. In the proof, we shall write β , etc, in place of $\beta^{(i)}$, etc. Taking δ sufficiently small, we start by choosing a small neighborhood W of S_j in $\bar{\Omega}_0$ globally parameterized by $(\sigma, y) \in [0, \delta] \times S_j$ (see the proof of Lemma A.5 in Appendix A). In fact, we can choose the coordinates and the small neighborhood of S_j such that $\sigma = cst$ corresponds to level sets for the function β (use $\nabla\beta$ for the vector field v in proof of Lemma A.5). Note that in the neighborhood W the function β decreases with σ .

Estimating from below the Jacobian¹ originating from the change of variable and observing that the integrand is constant w.r.t. y we obtain

$$\begin{aligned} W_{0,j}(s, \lambda) &\geq Cr^2|S_j| \int_0^\delta e^{-2s\eta(T', \sigma)} \varphi^{\frac{3}{2}}(T', \sigma) d\sigma \\ &= Cr^2|S_j| e^{-2s\eta(T', S)} \int_0^\delta e^{-2s(\eta(T', \sigma) - \eta(T', S))} \varphi^{\frac{3}{2}}(T', \sigma) d\sigma. \end{aligned}$$

We now use the change of variables $\sigma' = \eta(T', \sigma) - \eta(T', S) \geq 0$ which yields

$$W_{0,j}(s, \lambda) \geq Cr^2|S_j| \lambda^{-1} e^{-2s\eta(T', S)} \int_0^{\delta'} e^{-2s\sigma'} \varphi^{\frac{1}{2}}(T', \sigma) |\partial_\sigma \beta|^{-1} d\sigma',$$

where $\delta' = \eta(T', \delta) - \eta(T', S)$. We can find in W a positive lower bound for $(\partial_\sigma \beta)^{-1}$ independent of δ , i.e. the size of W . We thus obtain

$$\begin{aligned} W_{0,j}(s, \lambda) &\geq Cr^2|S_j| \lambda^{-1} e^{-2s\eta(T', S)} \varphi^{\frac{1}{2}}(T', \delta) \int_0^{\delta'} e^{-2s\sigma'} d\sigma' \\ &\geq Cr^2|S_j| s^{-1} \lambda^{-1} e^{-2s\eta(T', S)} \varphi^{\frac{1}{2}}(T', \delta) \int_0^{s\delta'} e^{-2\sigma'} d\sigma' \\ &\geq C'(s) r^2 |S_j| s^{-1} \lambda^{-1} e^{-2s\eta(T', S)} \varphi^{\frac{1}{2}}(T', \delta), \end{aligned}$$

with $C'(s)$ increasing with s . Observe now that

$$\begin{aligned} \varphi(T', \delta) &= \frac{e^{\lambda\beta(\delta)}}{(T - T')(T' - t_0)} = \frac{e^{\lambda\beta(S)}}{(T - T')(T' - t_0)} e^{\lambda(\beta(\delta) - \beta(S))} \\ &= \varphi(T', S) e^{\lambda(\beta(\delta) - \beta(S))}. \end{aligned}$$

Choosing δ sufficiently small such that $\frac{1}{2}(\beta(S) - \beta(\delta)) \leq \varepsilon$ thus yields the result. \blacksquare

With the previous lemmas we can now prove that the coefficient of $|\xi_{0,j}|^2$, $j = 1, \dots, p_0$, in (2.8) can be made positive. This requires taking *both* λ and s sufficiently large.

Proposition 2.7. *Let $1 \leq j \leq p_0$. There exists $\lambda_{2,j} \geq \lambda_1$ such that if $\lambda \geq \lambda_{2,j}$ then for s sufficiently large*

$$\begin{aligned} A_{0,j} &= (1 - CK'^2 s^{-\frac{3}{2}} \lambda^{-2}) W_{0,j}(s, \lambda) - CM' [s^{-\frac{1}{2}} \lambda^{-1} w_1(s, \lambda) \\ &\quad + s^{-\frac{3}{2}} \lambda^{-2} w_4(s, \lambda) + s^{-\frac{7}{2}} \lambda^{-2} w_0(s, \lambda)] \geq C(s, \lambda) > 0, \end{aligned}$$

with $C(s, \lambda) = C(s, \lambda, r, K, M, c_{max}, c_{min}, j)$.

¹Note that the estimation from below of the Jacobian is independent from the size of the neighborhood W .

Proof. It suffices to prove the result for $w_k^{(i)}(s, \lambda)$ and $W_{0,j}^{(i)}(s, \lambda)$. We shall write β , etc, in place of $\beta^{(i)}$, etc. We take s sufficiently large such that $(1 - CK'^2 s^{-\frac{3}{2}} \lambda^{-2}) \geq c_0 > 0$.

From Lemmas 2.5 and 2.6, for $\varepsilon > 0$ we obtain

$$(2.10) \quad A_{0,j} \geq -s^{-2} \nu(\lambda) C_1 e^{-2s\eta(T', S)} + \frac{1}{s\lambda} e^{-2s\eta(T', S)} \left[\frac{C_0}{2} C_\varepsilon r^2 |S_j| (\varphi(T', S))^{\frac{1}{2}} e^{-\lambda\varepsilon} - CM' \frac{\sqrt{\pi} \varphi(T', S)}{\sqrt{\phi''(T')} \sqrt{e^{\lambda\bar{\beta}} - e^{\lambda\beta(S)}}} \right],$$

where ν are a bounded function. We first treat the second term in the previous expression. Note that this term originates from the estimate from below for $W_{0,j}^{(i)}$ and the estimate of $s^{-\frac{1}{2}} \lambda^{-1} w_1^{(i)}$ by Lemma 2.5. The other terms in $A_{0,j}$ and the remainder part of the estimation of $s^{-\frac{1}{2}} \lambda^{-1} w_1^{(i)}$ are lumped in the first term of (2.10).

Choose now $\varepsilon < \frac{1}{2}(\bar{\beta} - \beta(S))$ (recall that $\bar{\beta} > \beta(S)$ because $m > 1$). Then since $\bar{\beta} > \beta(S)$ we have

$$\frac{\varphi(T', S)}{\sqrt{e^{\lambda\bar{\beta}} - e^{\lambda\beta(S)}}} = o((\varphi(T', S))^{\frac{1}{2}} e^{-\lambda\varepsilon})$$

for λ large. Thus the second term can be made positive for λ , say $\lambda = \lambda_{2,j}$, sufficiently large.

Once λ is fixed larger than $\lambda_{2,j}$, the first term in (2.10) can be made positive by taking s sufficiently large. \blacksquare

We now prove that the coefficient of $|\xi_{1,j}|^2$, $j = 1, \dots, p_1$, in (2.8) can be made positive. Here, the parameter λ is not of use.

Proposition 2.8. *Let $1 \leq j \leq p_1$. Let $\lambda \geq \lambda_1$. Then for s sufficiently large*

$$A_{1,j} = (1 - CK'^2 s^{-\frac{3}{2}} \lambda^{-2}) W_{1,j}(s, \lambda) - CM' [s^{-\frac{1}{2}} \lambda^{-1} w_1(s, \lambda) + s^{-\frac{3}{2}} \lambda^{-2} w_4(s, \lambda) + s^{-\frac{7}{2}} \lambda^{-2} w_0(s, \lambda)] \geq C(s, \lambda) > 0,$$

with $C(s, \lambda) = C(s, \lambda, r, K, M, c_{\max}, c_{\min})$.

Proof. It suffices to prove the result for $w_k^{(i)}(s, \lambda)$ and $W_{1,j}^{(i)}(s, \lambda)$. We shall write β , etc, in place of $\beta^{(i)}$, etc. We take s sufficiently large such that $(1 - CK'^2 s^{-\frac{3}{2}} \lambda^{-2}) \geq C_0 > 0$.

We first write

$$e^{-2s\eta(t, S)} = e^{-2s\eta(T', S)} e^{-2s(\eta(t, S) - \eta(T', S))},$$

and observe that for $s \geq s_0 > 0$

$$\int_{t_0}^T e^{-2s(\eta(t,S)-\eta(T',S))} \varphi^k(t,S) dt \leq L(s_0, \lambda, k),$$

for some positive $L(s_0, \lambda, k)$. From Lemma 1.1, there exists $V \Subset \Omega_{1,j}$ such that $\inf_{x \in V} \beta > \beta(S)$. Then with

$$\eta_{max}^{T',V} = \sup_{x \in V} \frac{e^{2\lambda K^{(i)}} - e^{\lambda \beta^{(i)}(x)}}{(T' - t_0)(T - T')}$$

we have $-\eta(T', x) \geq -\eta_{max}^{T',V} > -\eta(T', S)$, for $x \in V$, and $s > 0$. These observations yield

$$W_{1,j}(s) \geq r^2 \int_V e^{-2s\eta(T',x)} \varphi^{\frac{3}{2}}(T', x) dx \geq C(\lambda) r^2 |V| e^{-2s\eta_{max}^{T',V}},$$

and

$$w_k(s) \leq L(s_0, \lambda, k) e^{-2s\eta(T',S)},$$

which implies the result. ■

With (2.8) and Propositions 2.7 and 2.8, recalling that $v = u_t = \partial_t(y - \tilde{y})$, We have thus obtained the following stability result.

Theorem 2.9. *Let γ be a subset of the boundary Γ of an open set Ω of \mathbb{R}^n that satisfies Condition (GC), and γ satisfies Assumption 0.4. We assume that the diffusion coefficients c and \tilde{c} satisfy Assumptions 0.2 and 2.4 and $c_0 - c_1 \geq \Delta > 0$. Let y_0, \tilde{y}_0 in $L^2(\Omega)$ and let y, \tilde{y} be solutions to (2.1)–(2.2) satisfying Assumption 2.1. Then there exists a constant C*

$$C = C(\Omega, T, t_0, \gamma, S, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, M, K, r, c_{min}, c_{max}, \Delta)$$

such that

$$(2.11) \quad \sum_{i=0,1} \sum_{j=1}^{p_i} |c_{ij} - \tilde{c}_{ij}|^2 \leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 + C |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\Omega')}^2.$$

We shall see in Proposition 3.5, below, that we can achieve the regularity properties and estimates of Assumption 2.1.

Remark 2.10. Observe than in the statement of Theorem 2.9 the initial condition y_0 and \tilde{y}_0 need not be equal (see systems (2.1)–(2.2)).

Remark 2.11. If the position of the interface S is known, we can improve the result of Theorem 2.9 by locally observaing of the solutions y and \tilde{y} at time T' . Let $\omega = \omega_0 \cup \omega_1$, with ω_1 a neighborhood of γ in $\overline{\Omega}_1$ and ω_0 a neighborhood of S in $\overline{\Omega}_0$. We can in fact relax Assumption 2.1 with the following one: Let $r > 0$,

the solution \tilde{y} is such that $|\Delta\tilde{y}(T')| \geq r > 0$ in ω . Then, we obtain a stability estimate by solely observing y and \tilde{y} at time T' on ω in place of Ω

(2.12)

$$\sum_{i=0,1} \sum_{j=1}^{p_i} |c_{ij} - \tilde{c}_{ij}|^2 \leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 + C |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\omega)}^2.$$

We briefly sketch the proof, as it closely follows the exposition of the proof of estimate (2.11) in this section. It relies on the following lemma.

Lemma 2.12. *Let $s_3 > 0$ and $\lambda_3 > 0$. There exists $C > 0$ such that for all $s \geq s_3$, and $\lambda \geq \lambda_3$, we have*

$$\begin{aligned} \int_{\Omega_i} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}(T',x)} \varphi^{(2)\frac{3}{2}} \right) (T', x) dx \\ \leq C \int_{\omega_i} \left(e^{-2s\eta^{(1)}} \varphi^{(1)\frac{3}{2}} + e^{-2s\eta^{(2)}(T',x)} \varphi^{(2)\frac{3}{2}} \right) (T', x) dx \end{aligned}$$

Proof. Set $A = \sup_{\Omega_1 \setminus \omega_1} \beta$ and let $\varepsilon > 0$ be sufficiently small such that $B = A + \varepsilon < \sup_{\Omega_1} \beta$ and define $\tilde{\omega}_1 = \{x \in \Omega_1; A + \varepsilon < \beta(x)\}$. Then $\tilde{\omega}_1 \subset \omega_1$ and we set $K = \frac{|\Omega_1 \setminus \omega_1|}{|\omega_1|}$. Defining $\eta_B^{(i)}$ and $\varphi_B^{(i)}$ to be equal to $\eta^{(i)}$ and $\varphi^{(i)}$ at time T' and $\beta(x)$ replaced by B . Then for all $s \geq s_3$, and $\lambda \geq \lambda_3$, we have

$$\int_{\Omega_1} e^{-2s\eta_B^{(i)}} \varphi_B^{(i)\frac{3}{2}} dx \leq K \int_{\omega_1} e^{-2s\eta_B^{(i)}} \varphi_B^{(i)\frac{3}{2}} dx, \quad i = 1, 2.$$

We conclude by observing that, for $i = 1, 2$,

$$\begin{aligned} e^{-2s\eta_B^{(i)}} &\leq e^{-2s\eta^{(i)}(T',x)}, \quad \varphi_B^{(i)} \leq \varphi^{(i)}(T',x), \quad x \in \tilde{\omega}_1, \\ e^{-2s\eta^{(i)}(T',x)} &\leq e^{-2s\eta_B^{(i)}}, \quad \varphi^{(i)}(T',x) \leq \varphi_B^{(i)}, \quad x \in \Omega_1 \setminus \tilde{\omega}_1. \end{aligned}$$

The same method can be applied for the second set of integrals on Ω_0 and ω_0 . \blacksquare

Continuation of Remark 2.11. In the statement of Lemma 2.3 we can replace the integration domain, Ω' , in the l.h.s. of the estimate, by ω .

To reach an equation of the form of (2.8) with the volume integrals computed² over ω in place of Ω we write the following counterpart to (2.7):

$$\begin{aligned} \sum_{i=0,1} \int_{t_0}^T \int_{\Omega_i} (e^{-2s\eta^{(1)}} + e^{-2s\eta^{(2)}}) |\xi \Delta \partial_t \tilde{y}|^2 dx dt \\ \leq K'^2 \sum_{i=0,1} \sum_{j=1}^{p_i} |\xi_{i,j}|^2 \int_{\Omega_{i,j} \cap \omega_i} \left(e^{-2s\eta^{(1)}(T',x)} + e^{-2s\eta^{(2)}(T',x)} \right) |\Delta \tilde{y}(T', x)|^2 dx \end{aligned}$$

by Lemma 2.12. The result of Lemma 2.6 remain unchanged as ω_0 is a neighborhood of S in $\bar{\Omega}_0$. In the proof of Proposition 2.8 we can choose the open set V to be in ω_1 .

²including the definition of $W_{i,j}$, $i = 0, 1$, $j = 1, \dots, p_i$.

Remark 2.13. Note that if we assume that $y(T', \cdot) = \tilde{y}(T', \cdot)$ then the stability estimate becomes

$$\sum_{i=0,1} \sum_{j=1}^{p_i} |c_{ij} - \tilde{c}_{ij}|^2 \leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2.$$

Such an additional assumption is sometimes made, e.g. in [19].

With Theorem 2.9 we have the following uniqueness result

Corollary 2.14. *Under the same assumptions as Theorem 2.9 and if*

$$\begin{aligned} \partial_n(\partial_t(y - \tilde{y}))(t, x) &= 0 \text{ in } (t_0, T) \times \gamma, \\ \Delta y(T', x) - \Delta \tilde{y}(T', x) &= 0, \text{ in } \Omega', \end{aligned}$$

then $c = \tilde{c}$. Furthermore $y_0 = \tilde{y}_0$.

Proof. The second assertion remains to be proved. If $c = \tilde{c}$ then $u = y - \tilde{y} \in \mathcal{D}_A$, with $A = \nabla \cdot (c \nabla(\cdot))$ (see Appendix A), is solution to

$$\begin{cases} \partial_t u - \nabla \cdot (c \nabla u) = 0 & \text{in } (0, T) \times \Omega, \\ u = 0, & \text{on } (0, T) \times \Gamma, \\ u(0, x) = u_0(x), & \text{in } \Omega, \end{cases}$$

with $u_0 = y_0 - \tilde{y}_0$. Thus $u = S(t)u_0$. We have $\Delta(u)(T') = 0$ in Ω' . Thus $\nabla \cdot (c \nabla u)|_{\Omega'}(T') = 0$. Since $u(T') \in \mathcal{D}_A$ we have $u(T') = 0$. Since the semigroup $S(t)$ generated by $-\nabla \cdot (c \nabla(\cdot))$ is analytic by Proposition A.3, we obtain that $S(t)u_0 = 0$ for all $t > 0$. The continuity in $t = 0^+$ yields $u_0 = 0$. ■

If we make further assumptions on the initial conditions y_0 and \tilde{y}_0 we can in fact obtain a stability result for these initial conditions as well. This is the subject of Section 4.

Remark 2.15. In the stability result obtained here, we have made the choice to make some of the measurements on part of the boundary $(0, T) \times \Gamma$. Derivation of a Carleman estimate, as in [7], with a r.h.s. with an ‘observation’ in an inner volume $(0, T) \times \omega$ of $(0, T) \times \Omega_1$ would yield a stability estimate like (2.11) with $|\partial_t y - \partial_t \tilde{y}|_{L^2((t_0, T) \times \omega)}$ in the r.h.s..

Remark 2.16. Observe that in place of Assumption 2.4, we could have solely assumed that the difference $\xi = c - \tilde{c}$ is piecewise constant. Then, we would replace $c \Delta u$ by $\nabla \cdot (c \nabla u)$ in the r.h.s. of (2.8). This would yield a stability estimate of the form

$$\begin{aligned} (2.13) \quad \sum_{i=0,1} \sum_{j=1}^{p_i} |c_{ij} - \tilde{c}_{ij}|^2 &\leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 + C |\nabla y(T', \cdot) - \nabla \tilde{y}(T', \cdot)|_{(L^2(\Omega'))^n}^2 \\ &\quad + C |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\Omega')}^2. \end{aligned}$$

In Remark 3.7, in Section 3, we show that Assumption 2.1 can be fulfilled, in the case where ξ is piecewise constant, if $|\nabla \tilde{c}|$ is sufficiently small and \tilde{c} sufficiently regular.

Remark 2.17. Here we have set the r.h.s. of the parabolic equations in (2.1) and (2.2) to zero. With a non-vanishing r.h.s., say $f(t, x)$, the same stability result holds once the hypothesis of Assumption 2.1 are fulfilled, since (2.3) is preserved. Assuming that the function $f(t, x)$ is bounded and sufficiently regular, say of class \mathcal{C}^k , with k sufficiently large, we can obtain the same results as in Proposition 3.5 of Section 3, by a slight modification of the argumentation. In this case also, the class of solutions y and \tilde{y} satisfying Assumption 2.1 is non empty.

3 Existence of solutions y, \tilde{y} satisfying Assumption 2.1

We propose a possible choice of boundary condition h and of initial condition \tilde{y}_0 to achieve the particular properties for the solutions y and \tilde{y} listed in Assumption 2.1 needed in the proof of Theorem 2.9 in Section 2.

We shall denote $S(t)$ (resp. $\tilde{S}(t)$) the analytic semi-group generated by unbounded operator A (resp. \tilde{A}) formally defined by $-\nabla \cdot (c\nabla(\cdot))$ (resp. $-\nabla \cdot (\tilde{c}\nabla(\cdot))$) on $L^2(\Omega)$ with domain (See appendix A)

$$\begin{aligned} \mathcal{D}_A &= \{u \in H_0^1(\Omega); \nabla \cdot (c\nabla u) \in L^2(\Omega)\}, \\ (\text{resp. } \mathcal{D}_{\tilde{A}} &= \{u \in H_0^1(\Omega); \nabla \cdot (\tilde{c}\nabla u) \in L^2(\Omega)\}). \end{aligned}$$

The convention we use here is ' $S(t) = e^{-tA}$ '.

Lemma 3.1. *Let $r > 0$ and let $\tilde{c} \in L^\infty(\Omega)$. There exists $\tilde{y}_0 \in \mathcal{D}_{\tilde{A}}$ and $\chi : [0, T] \rightarrow \mathbb{R}$ such that the solution to*

$$(3.1) \quad \begin{cases} \partial_t \tilde{y} - \nabla \cdot (\tilde{c}\nabla \tilde{y}) = 0, & \text{in } (0, T) \times \Omega, \\ \tilde{y}(t, x) = \chi(t), & \text{on } (0, T) \times \Gamma, \\ \tilde{y}(t, \cdot) - \chi(t) \in \mathcal{D}_{\tilde{A}}, & 0 < t \leq T, \\ \tilde{y}(0) = \tilde{y}_0, \end{cases}$$

satisfies $|\nabla \cdot (\tilde{c}\nabla \tilde{y})(T')| \geq r > 0$ a.e.. The function χ can be chosen such that χ' is a positive constant.

Proof. Observe that $p(t, x) = \tilde{y}(t, x) - \chi(t)$ is solution to

$$(3.2) \quad \begin{cases} \partial_t p - \nabla \cdot (\tilde{c}\nabla p) = -\chi'(t), & \text{in } (0, T) \times \Omega, \\ p(t, x) = 0, & \text{on } (0, T) \times \Gamma, \\ p(t, \cdot) \in \mathcal{D}_{\tilde{A}}, & 0 < t \leq T, \\ p(0, x) = \tilde{y}_0(x) - \chi(0) = p_0 \in L^2(\Omega), \end{cases}$$

and is thus given by Duhamel's formula [21]

$$(3.3) \quad p(t) = \tilde{S}(t)p_0 - \int_0^t \tilde{S}(t-s)\chi'(s)ds.$$

In fact, we choose χ of the form $\chi(t) = -\rho t$, where ρ is a negative constant. We also choose \tilde{y}_0 such that $p_0 = \tilde{y}_0 \in \mathcal{D}_{\tilde{A}}$ and $\nabla \cdot (\tilde{c} \nabla \tilde{y}_0) \geq r_0 > r$ a.e. in Ω (choose $f \in L^2(\Omega)$, such that $f > r_0$, and solve the elliptic problem $\nabla \cdot (\tilde{c} \nabla \tilde{y}_0) = f$ for \tilde{y}_0 in $H_0^1(\Omega)$). We choose ρ such that $-r_0 < \rho \leq -r < 0$.

The solution p to (3.2) is unique in $\mathcal{C}^1([0, T], L^2(\Omega)) \cap \mathcal{C}^0([0, T], \mathcal{D}_{\tilde{A}})$ and given by (3.3) [6, Theorem 3 and following Remark 2, Section XVII B.1]. Denoting by $\mathbf{1}$ the function identically equal to 1 on Ω we find

$$p(t) = \tilde{S}(t)p_0 + \rho \int_0^t \tilde{S}(s)\mathbf{1}ds,$$

which yields $q := -\tilde{A}p + \rho\mathbf{1} := \tilde{S}(t)(\nabla \cdot (\tilde{c} \nabla p_0) + \rho\mathbf{1})$ [21, Theorem 1.2.4]. Hence q is the solution to

$$\begin{cases} \partial_t q - \nabla \cdot (\tilde{c} \nabla q) = 0, & \text{in } (0, T) \times \Omega, \\ q(t, x) = 0, & \text{on } (0, T) \times \Gamma, \\ q(t, \cdot) \in \mathcal{D}_{\tilde{A}}, & 0 < t \leq T, \\ q(0, x) = q_0(x) := \nabla \cdot (\tilde{c} \nabla p_0) + \rho\mathbf{1}. \end{cases}$$

We now apply the maximum principle (which is valid for L^∞ diffusion coefficients) [5, proof of Theorem IX.3] which reads for the time interval $[0, T']$

$$\text{ess inf}_{\bar{Q}_{T'}} q \geq \min(0, \text{ess inf}_{\Omega} q_0) = 0, \quad Q_{T'} = (0, T') \times \Omega.$$

This yields $\nabla \cdot (\tilde{c} \nabla p)(T', x) \geq -\rho \geq r > 0$ a.e.. ■

Lemma 3.2. *Let $l > n/2$ and $l \geq 2$. Let $\tilde{c} \in L^\infty(\Omega)$ be such that $\tilde{c}|_{\Omega_i}$ is $\mathcal{C}^{l-1}(\bar{\Omega}_i)$, $i = 0, 1$. Let Ω be such that S and $\partial\Omega$ are of class \mathcal{C}^l . Let $\tilde{y}_0 \in \mathcal{D}_{\tilde{A}}$ and the function $\chi : [0, T] \rightarrow \mathbb{R}$, such that χ' is constant, be both chosen according to Lemma 3.1. Then $\nabla \cdot (\tilde{c} \nabla \tilde{y})|_{\Omega_i} \in \mathcal{C}^k((0, T], L^\infty(\Omega_i))$, $i = 0, 1$, for all $k \in \mathbb{N}$. Let $\varepsilon > 0$ then $\nabla \cdot (\tilde{c} \nabla \tilde{y})|_{\Omega_i}$, $i = 0, 1$, remain in a bounded domain of $\mathcal{C}^k([\varepsilon, T], L^\infty(\Omega_i))$ for all $k \in \mathbb{N}$, uniformly w.r.t. \tilde{c} and \tilde{y}_0 , for $0 < c_{\min} \leq \tilde{c} \leq c_{\max}$ and $\nabla \cdot (\tilde{c} \nabla \tilde{y}_0)$ in a bounded domain of $L^2(\Omega)$.*

Proof. We use the notations of the proof of Lemma 3.1. We set $p(t, x) = \tilde{y}(t, x) - \chi(t)$ and observe that $q := -\tilde{A}p + \rho\mathbf{1}$ is the solution to

$$\begin{cases} \partial_t q - \nabla \cdot (\tilde{c} \nabla q) = 0, & \text{in } (0, T) \times \Omega, \\ q(t, x) = 0, & \text{on } (0, T) \times \Gamma, \\ q(t, \cdot) \in \mathcal{D}_{\tilde{A}}, & 0 < t \leq T, \\ q(0, x) = q_0(x) := \nabla \cdot (\tilde{c} \nabla p_0) + \rho\mathbf{1}. \end{cases}$$

From Corollary A.11 we have that $q|_{(0, T] \times \Omega_i} \in \mathcal{C}^k((0, T]; H^l(\Omega_i))$, $i = 0, 1$, for all $k \in \mathbb{N}$. Since $l > n/2$, the space $H^l(\Omega_i)$ is continuously embedded in $L^\infty(\Omega_i)$ which yields the result. The last statement follows from Remark A.12. ■

Remark 3.3. In the case of $n=2, 3$, which concerns most of the applications, we choose $m = 2$. The condition on S , $\partial\Omega$ and the coefficients $\tilde{c}|_{\Omega_i}$ in the previous lemma are then the default ones assumed in the introduction.

Let the function χ , such that $\rho = -\chi'$ is constant, be chosen according to Lemma 3.1. We then have the following regularity property.

Lemma 3.4. *Let $c, \tilde{c} \in L^\infty(\Omega)$ be such that $c|_{\Omega_i}, \tilde{c}|_{\Omega_i}$ is $\mathcal{C}^1(\overline{\Omega}_i)$, $i = 0, 1$, $0 < c_{\min} \leq c, \tilde{c} \leq c_{\max}$ and $y_0, \tilde{y}_0 \in L^2(\Omega)$ remain in a bounded domain of $L^2(\Omega)$. The solutions \tilde{y} and y to*

$$\begin{cases} \partial_t \tilde{y} - \nabla \cdot (\tilde{c} \nabla \tilde{y}) = 0, & \text{in } (0, T) \times \Omega, \\ \tilde{y}(t, x) = \chi(t), & \text{on } (0, T) \times \Gamma, \\ \text{trans. condition (TC1),} \\ \tilde{y}(0) = \tilde{y}_0, \end{cases} \quad \begin{cases} \partial_t y - \nabla \cdot (c \nabla y) = 0, & \text{in } (0, T) \times \Omega, \\ y(t, x) = \chi(t), & \text{on } (0, T) \times \Gamma, \\ \text{trans. condition (TC1),} \\ y(0) = y_0, \end{cases}$$

belong to $\mathcal{C}^k((0, T], H^1(\Omega))$ and are such that $\tilde{y}|_{\Omega_i}, y|_{\Omega_i} \in \mathcal{C}^k((0, T], H^2(\Omega_i))$, $i = 0, 1$, for all $k \in \mathbb{N}$. Let $\varepsilon > 0$, then for all $k \in \mathbb{N}$, $\tilde{y}|_{\Omega_i}, y|_{\Omega_i}$ remain in a bounded domain of $\mathcal{C}^k((\varepsilon, T], H^2(\Omega_i))$, $i=0,1$, uniformly w.r.t. \tilde{c} , y_0 , and \tilde{y}_0 .

Proof. We work out the proof for y . We define $p(t, x) = y(t, x) - \chi(t)$. The function p is solution to

$$(3.4) \quad \begin{cases} \partial_t p - \nabla \cdot (c \nabla p) = \rho, & \text{in } (0, T) \times \Omega, \\ p(t, x) = 0, & \text{on } (0, T) \times \Gamma, \\ p(t, \cdot) \in \mathcal{D}_A, & 0 < t \leq T, \\ p(0) = p_0 = y_0 - \chi(0) \in L^2(\Omega). \end{cases}$$

It suffices to prove the result for p . Since ρ is constant, the (mild) solution to (3.4) is a classical solution [21, Theorem 4.3.2]. We prove below that $p \in \mathcal{C}^k((0, T], \mathcal{D}_A)$, $k > 0$. Thus $p \in \mathcal{C}^k((0, T], L^2(\Omega))$. Since $D_A \subset H_0^1$ with continuous injection then $p \in \mathcal{C}^k((0, T], H_0^1(\Omega))$. By Proposition A.4 the maps $p \mapsto p|_{\Omega_i}$, $i = 0, 1$, are continuous from \mathcal{D}_A into $H^2(\Omega_i)$. Thus $p|_{\Omega_i} \in \mathcal{C}^k((0, T], H^2(\Omega_i))$.

The solution p is given by

$$p(t) = S(t)p_0 + \rho \int_0^t S(s) \mathbf{1} ds,$$

where $\mathbf{1}$ is the function identically equal to 1 on Ω . The first term $p_1 = S(t)p_0$ in $\mathcal{C}^k((0, T], \mathcal{D}_A)$ for all $k, l > 0$ by Proposition A.2. For the second term $p_2 = \rho \int_0^t S(s) \mathbf{1} ds$ we have [21, Theorem 1.2.4] $-Ap_2 = \rho(S(t)\mathbf{1} - \mathbf{1})$. Thus $Ap_2 \in \mathcal{C}^k((0, T], L^2(\Omega))$, i.e. $p_2 \in \mathcal{C}^k((0, T], \mathcal{D}_A)$, for all $k > 0$. The boundedness statement follows from Remark A.12. \blacksquare

With the proposed initial condition \tilde{y}_0 and boundary condition $h(t, x) = \chi(t)$ we have thus obtained the following regularity and boundedness properties.

Proposition 3.5. *Let $r > 0$. Let $y_0 \in L^2(\Omega)$. Let Ω be such that S and $\Gamma = \partial\Omega$ are of class \mathcal{C}^l and $\tilde{c} \in L^\infty(\Omega)$ be such that $\tilde{c}|_{\Omega_i}$ is $\mathcal{C}^{l-1}(\overline{\Omega}_i)$, $i = 0, 1$, with $l > n/2$, $l \geq 2$. There exists $h(t, x) \in \mathcal{C}([0, T] \times \Gamma)$ and an initial condition $\tilde{y}_0 \in \mathcal{D}_{\tilde{A}}$ such that the solutions y, \tilde{y} to systems (2.1)–(2.2) satisfy*

1. $\nabla \cdot (\tilde{c} \nabla \tilde{y})(T') \geq r > 0$;
2. $\nabla \cdot (\tilde{c} \nabla \tilde{y})|_{\Omega_i} \in \mathcal{C}^k([t_0, T], L^\infty(\Omega_i))$, $i = 0, 1$, for all $k \in \mathbb{N}$;
3. $y, \tilde{y} \in \mathcal{C}^k([t_0, T], L^2(\Omega)) \cap \mathcal{C}^k((t_0, T], H^1(\Omega))$, for all $k \in \mathbb{N}$;
4. $y|_{\Omega_i}, \tilde{y}|_{\Omega_i} \in \mathcal{C}^k([t_0, T], H^2(\Omega_i))$, $i = 0, 1$, for all $k \in \mathbb{N}$.

The restrictions $\tilde{y}|_{\Omega_i}$ remain in a bounded domain of $\mathcal{C}^k([t_0, T], H^2(\Omega_i))$ and $\nabla \cdot (\tilde{c} \nabla \tilde{y})|_{\Omega_i}$ in a bounded domain of $\mathcal{C}^k([t_0, T], L^\infty(\Omega_i))$, uniformly w.r.t. \tilde{c} and \tilde{y}_0 if $0 < c_{\min} \leq \tilde{c} \leq c_{\max}$ and $\nabla \cdot (\tilde{c} \nabla \tilde{y}_0)$ remain in a bounded domain of $L^2(\Omega)$.

With Proposition 3.5 we observe that Assumption 2.1 in Section 2 can be fulfilled in the framework of Assumption 2.4 when $c_{\min} \leq \tilde{c} \leq c_{\max}$ and $\tilde{y}_0 \in \mathcal{D}_{\tilde{A}}$ such that $\nabla \cdot (\tilde{c} \nabla \tilde{y}_0)$ remain in a bounded domain of $L^2(\Omega)$ for properly chosen boundary conditions $h(t, x)$.

Remark 3.6. Observe that we could simply assume that $\tilde{y}_0 \in L^2(\Omega)$ and design the boundary condition $h(t, x)$ to reach a proper state in $\mathcal{D}_{\tilde{A}}$ in a finite time $t_1 < t_0$. This can be achieved as the parabolic equation we study here is null-controllable, i.e. exactly controllable to the trajectories [7].

Remark 3.7. In the case where we only assume that ξ is piecewise constant, then from Item 1 of Proposition 3.5, we can obtain Item 1 of Assumption 2.1, in the case where $|\nabla \tilde{c}|$ is sufficiently small, since, from the proof of Lemma 3.2, we also find that $\nabla \tilde{y}|_{\Omega_i}$, $i = 0, 1$, remain in bounded domains of $\mathcal{C}^k([\varepsilon, T], (L^\infty(\Omega_i))^n)$. Note that we need to assume a sufficiently large regularity on the coefficient \tilde{c} in the proof of Lemma 3.2.

4 Uniqueness and stability estimate for the initial conditions

In this section we closely follow the method of [22]. We shall assume

Assumption 4.1. Let $r_0 > 0$. The initial conditions y_0 and \tilde{y}_0 satisfy

1. y_0 is in a bounded domain of \mathcal{D}_A ;
2. \tilde{y}_0 is in a bounded domain of $\mathcal{D}_{\tilde{A}}$;
3. $\nabla \cdot (\tilde{c} \nabla \tilde{y}_0) \geq r_0$,
4. y, \tilde{y} are in a bounded domain of $\mathcal{C}^1([0, T], L^2(\Omega))$, where y and \tilde{y} are the solutions to (2.1)–(2.2).

Observe that item 4 of Assumption 4.1 implies that $\tilde{y}|_{\Omega_i}, y|_{\Omega_i}$ are in a bounded domain of $\mathcal{C}([0, T], H^2(\Omega_i))$, $i = 0, 1$ by Proposition A.4.

We denote $\tilde{z} = \partial_t \tilde{y} \in \mathcal{C}([0, T], L^2(\Omega))$ and thus $\tilde{z}(0)$ is well defined in $L^2(\Omega)$. We introduce w the solution to

$$(4.1) \quad \begin{cases} \partial_t w - \nabla \cdot (c \nabla w) = 0, & \text{in } (0, T) \times \Omega, \\ w(t, x) = \partial_t h(t, x), & \text{on } (0, T) \times \Gamma, \\ \text{transmission conditions (TC1),} & \text{on } (0, T) \times S, \\ w(0) = \tilde{z}(0), \end{cases}$$

and we further assume

Assumption 4.2. *The functions \tilde{z}, w are in a bounded domain of $L^2(0, T, H^1(\Omega))$.*

We also assume that the diffusion coefficients c and \tilde{c} are piecewise constant (Assumption 2.4).

Observe that if we choose the boundary condition $h(t, x) = -\rho t$ for $0 < t \leq T$ according to the proof of Lemma 3.1 (with $0 < r < r_0$), then the above assumption are fulfilled. In fact, the results of Section 3 show that Assumption 2.1 is then satisfied. In addition, item 4 in Assumption 4.1 and Assumption 4.2 are fulfilled by the following lemma.

Lemma 4.3. *If $h(t, x) = -\rho t$ then the solutions y, \tilde{y} to (2.1)–(2.2) and w to (4.1) satisfy item 4 in Assumption 4.1 and Assumption 4.2.*

Proof. We prove $w \in L^2(0, T, H^1(\Omega))$. The proof is the same for \tilde{z} . Let $p(t, x) = w(t, x) - \partial_t h(t, x) = w(t, x) + \rho$. Then p satisfies

$$\begin{cases} \partial_t p - \nabla \cdot (c \nabla p) = 0, & \text{in } (0, T) \times \Omega, \\ p(t, x) = 0, & \text{on } (0, T) \times \Gamma, \\ \text{transmission conditions (TC1),} & \text{on } (0, T) \times S, \\ p(0) = \tilde{z}(0) + \rho \in L^2(\Omega), \end{cases}$$

We thus have the usual energy estimate

$$\frac{1}{2} |p(t)|_{L^2(\Omega)}^2 + \int_0^t \int_{\Omega} c |\nabla p|^2 dt dx = \frac{1}{2} |p(0)|_{L^2(\Omega)}^2,$$

and p , and thus w , is in $L^2(0, T, H^1(\Omega))$ and remains in a bounded domain of this space if $c_{\min} \leq c \leq c_{\max}$ and \tilde{y}_0 remains in a bounded domain of $\mathcal{D}_{\tilde{A}}$.

To prove that y is in a bounded domain of $\mathcal{C}^1([0, T], L^2(\Omega))$ (the proof is the same for \tilde{y}), we set $p(t, x) = y(t, x) + \rho t$ and observe that $q := -Ap + \rho \mathbf{1}$ is $\mathcal{C}([0, T], L^2(\Omega))$ and thus $p \in \mathcal{C}([0, T], \mathcal{D}_A)$. Then $p \in \mathcal{C}^1([0, T], L^2(\Omega))$. ■

Define v_1 and v_2 that satisfy

$$(4.2) \quad \begin{cases} \partial_t v_1 - \nabla \cdot (c \nabla v_1) = \nabla \cdot (\xi \nabla \partial_t \tilde{y}), & \text{in } (0, T) \times \Omega', \\ v_1 = 0, & \text{on } (0, T) \times \Gamma, \\ \text{trans. conditions (TC2),} & \text{on } (0, T) \times S, \\ v_1(0) = 0, \end{cases}$$

and

$$(4.3) \quad \begin{cases} \partial_t v_2 - \nabla \cdot (c \nabla v_2) = 0, & \text{in } (0, T) \times \Omega', \\ v_2 = 0, & \text{on } (0, T) \times \Gamma, \\ \text{trans. conditions (TC1),} & \text{on } (0, T) \times S, \\ v_2(0) = \partial_t(y - \tilde{y})(0). \end{cases}$$

Observe that $\partial_t(y - \tilde{y})(0)$ is well defined and in a bounded domain of $L^2(\Omega)$ by Assupmtion 4.1.

With an argument of logarithmic convexity we have

$$(4.4) \quad |v_2(t)|_{L^2(\Omega)} \leq K^{1-t/T'} |v_2(T')|_{L^2(\Omega)}^{t/T'}, \quad 0 \leq t \leq T'.$$

Such an estimate makes use of the convexity of $F(t) = \ln(|v_2(t)|_{L^2(\Omega)}^2)$ and $|v_2(0)|_{L^2(\Omega)} \leq K$ (see [20, Section 2.3] for further details).

We now prove the following lemma

Lemma 4.4. *There exists $C > 0$, such that*

$$|v_1(t)|_{L^2(\Omega)} \leq C |c - \tilde{c}|_{L^\infty(\Omega)}^{\frac{1}{2}}, \quad 0 \leq t \leq T.$$

Note that v_1 satisfies transmission condition (TC2) and thus does not belong to \mathcal{D}_A . We thus cannot use some argument of regularity w.r.t. the source term from parabolic theory for v_1 .

Proof. First observe that $v_1 = w - \tilde{z}$. From Assumption 4.2, w and \tilde{z} remain in a bounded domain of $L^2(0, T, H^1(\Omega))$. Now $\partial_t(w - \tilde{z}) - \nabla \cdot (c \nabla w - \tilde{c} \nabla \tilde{z}) = 0$, which after multiplication by $w - \tilde{z}$, integration over Ω and an integration by parts, yields

$$\begin{aligned} 0 &= \frac{1}{2} \partial_t \int_{\Omega} |w - \tilde{z}|^2 dx + \int_{\Omega} (c \nabla w - \tilde{c} \nabla \tilde{z}) \cdot (\nabla w - \nabla \tilde{z}) dx \\ &= \frac{1}{2} \partial_t \int_{\Omega} |w - \tilde{z}|^2 dx + \int_{\Omega} (c - \tilde{c}) \nabla w \cdot (\nabla w - \nabla \tilde{z}) dx + \int_{\Omega} \tilde{c} |\nabla w - \nabla \tilde{z}|^2 dx \end{aligned}$$

since $\nabla \cdot (c \nabla w - \tilde{c} \nabla \tilde{z}) \in L^2(\Omega)$ by the definitions of \mathcal{D}_A and $\mathcal{D}_{\tilde{A}}$ (see Appendix A). We thus obtain

$$\begin{aligned} \frac{1}{2} \partial_t \int_{\Omega} |w - \tilde{z}|^2 dx &\leq \left| \int_{\Omega} (c - \tilde{c}) \nabla w \cdot (\nabla w - \nabla \tilde{z}) dx \right| \\ &\leq |c - \tilde{c}|_{L^\infty(\Omega)} |\nabla w|_{L^2(\Omega)} |\nabla w - \nabla \tilde{z}|_{L^2(\Omega)}. \end{aligned}$$

Integrating over $(0, t)$ yields the result. \blacksquare

As in Section 2, we define $v = \partial_t(y - \tilde{y})$ and observe that $v = v_1 + v_2$. We thus have

$$\begin{aligned} |v(t)|_{L^2(\Omega)} &\leq |v_1(t)|_{L^2(\Omega)} + |v_2(t)|_{L^2(\Omega)} \leq C |c - \tilde{c}|_{L^\infty(\Omega)}^{\frac{1}{2}} + |v_2(T')|_{L^2(\Omega)}^{t/T'} \\ &\quad 0 \leq t \leq T', \end{aligned}$$

and

$$\begin{aligned} |v_2(T')|_{L^2(\Omega)} &\leq |v(T')|_{L^2(\Omega)} + |v_1(T')|_{L^2(\Omega)} \\ &\leq C(|(\Delta y - \Delta \tilde{y})(T')|_{L^2(\Omega)} + |c - \tilde{c}|_{L^\infty(\Omega)} + |c - \tilde{c}|_{L^\infty(\Omega)}^{\frac{1}{2}}), \end{aligned}$$

making use of

$$v = \nabla \cdot (c \nabla u) + \nabla(\xi \cdot \nabla \tilde{y}) = c(\Delta y - \Delta \tilde{y}) + \xi \Delta \tilde{y}, \quad \text{in } \Omega',$$

where Assumption 2.4 was applied. This yields

$$|v(t)|_{L^2(\Omega)} \leq C \left(|c - \tilde{c}|_{L^\infty(\Omega)}^{\frac{1}{2}} + \nu^{t/T'} \right), \quad 0 \leq t \leq T',$$

with $\nu = |(y - \tilde{y})(T')|_{H^2(\Omega')} + |c - \tilde{c}|_{L^\infty(\Omega)} + |c - \tilde{c}|_{L^\infty(\Omega)}^{\frac{1}{2}}$. Because of the regularity of v w.r.t. time t we now have

$$\begin{aligned} |y_0 - \tilde{y}_0|_{L^2(\Omega)} &= |u_0|_{L^2(\Omega)} = \left| \int_0^{T'} v(t) dt - u(T') \right|_{L^2(\Omega)} \\ &\leq \int_0^{T'} |v(t)|_{L^2(\Omega)} dt + |y(T') - \tilde{y}(T')|_{L^2(\Omega)}, \end{aligned}$$

which gives

$$|y_0 - \tilde{y}_0|_{L^2(\Omega)} \leq C \left(T'^{\frac{\nu-1}{\ln(\nu)}} + T' |c - \tilde{c}|_{L^\infty(\Omega)}^{\frac{1}{2}} \right) + |y(T') - \tilde{y}(T')|_{L^2(\Omega)}.$$

We thus have

$$|y_0 - \tilde{y}_0|_{L^2(\Omega)} \leq C \left(\frac{\nu-1}{\ln(\nu)} + C' \nu \right).$$

Observing that $x < \frac{x-1}{\ln(x)}$ for $x \in (0, 1)$ we obtain

$$|y_0 - \tilde{y}_0|_{L^2(\Omega)} \leq C \frac{\nu-1}{\ln(\nu)}, \quad \text{if } \nu < 1.$$

With Theorem 2.9 we obtain

$$(4.5) \quad |c - \tilde{c}|_{L^\infty(\Omega)}^2 \leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 + C |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\Omega')}^2.$$

When the r.h.s. of (4.5) is sufficiently small, we obtain

$$0 < \nu \leq C \left(|\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 + |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\Omega')}^2 \right)^{\frac{1}{4}} < 1.$$

In that case

$$\ln(\nu) \leq C' \ln(C |(y - \tilde{y})(T')|_{H^2(\Omega')}^2 + C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2) < 0,$$

and thus

$$\frac{\nu-1}{\ln(\nu)} \leq C'(\nu-1) / \ln(C |(y - \tilde{y})(T')|_{H^2(\Omega')}^2 + C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2).$$

We thus obtain the following stability theorem for the initial conditions.

Theorem 4.5. *Under the hypothesis of Theorem 2.9 in addition to Assumptions 4.1 and 4.2 there exist some constants $C, C' > 0$*

$$C = C(\Omega, T, t_0, \gamma, S, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, M, K, r, c_{\min}, c_{\max}, \Delta),$$

$$C' = C'(\Omega, T, t_0, \gamma, S, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, M, K, r, c_{\min}, c_{\max}, \Delta)$$

such that

$$|y_0 - \tilde{y}_0|_{L^2(\Omega)} \leq C' / |\ln(C|(y - \tilde{y})(T')|_{H^2(\Omega')} + C|\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)})|$$

for $|(y - \tilde{y})(T')|_{H^2(\Omega')} + |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}$ sufficiently small.

5 A generalization to embedded materials

Let $\Omega \subset \mathbb{R}^n$ be a bounded connected open set. As before $\bar{\Omega}$ is assumed to be a \mathcal{C}^2 submanifold with boundary in \mathbb{R}^n . We consider some open sets V_0, \dots, V_m such that

$$V_0 \Subset V_1 \Subset \dots \Subset V_m = \Omega.$$

We then set $\Omega_i := V_i \setminus \bar{V}_{i-1}$, $i = 1, \dots, m$ and $\Omega_0 = V_0$. Each open set Ω_i , $i = 0, \dots, m$, is assumed to have p_i connected components, $\Omega_{i,j}$, $j = 1, \dots, p_i$. We let $S_{i-1}^i = \bar{\Omega}_{i-1} \cap \bar{\Omega}_i$ be the interface between Ω_{i-1} and Ω_i , $i = 1, \dots, m$. The interface S_{i-1}^i is assumed to be \mathcal{C}^2 . Similarly, we denote by $S_{i-1,j}^{i,k}$ the interface (possibly empty) between $\Omega_{i-1,j}$ and $\Omega_{i,k}$. We also set

$$S_{i,j} = \bigcup_{k=1, \dots, p_{i+1}} S_{i,j}^{i+1,k}, \text{ for } i = 0, \dots, m-1,$$

$$S^{i,j} = \bigcup_{k=1, \dots, p_{i-1}} S_{i-1,k}^{i,j}, \text{ for } i = 1, \dots, m,$$

and $S = \bigcup_{i=0, \dots, m-1} S_i^{i+1}$. On an interface S_{i-1}^i , we denote by n the outward unit normal vector to Ω_i and the outward unit normal vector to Ω_m on Γ . We set $\Omega' = \Omega_0 \cup \dots \cup \Omega_m$.

We introduce the following geometric condition.

Assumption 5.1. Geometric Condition (GC_m)

Let $\Omega_{i,j}$ be a component of Ω_i . We assume that either

Case 1 $i = 0$ or $\Omega_{i,j}$ has no interface in common with any component of Ω_{i-1} , i.e. $S^{i,j} = \emptyset$.

Case 2 $\Omega_{i,j}$ has an interface with only one component of Ω_{i-1} .

In **case 1** we further assume that there exists two disjoint open subsets $\mathcal{O}_{i,j}^{(1)}, \mathcal{O}_{i,j}^{(2)} \Subset \Omega_{i,j}$ and two vector fields, $\zeta_{i,j}^{(k)} \in \mathcal{C}^1(\bar{\Omega}_{i,j}, \mathbb{R}^2)$, $k = 1, 2$, such that

$$\zeta_{i,j}^{(k)}(x) \cdot n(x) > 0, \quad \forall x \in S_{i,j}, \quad k = 1, 2,$$

$$\zeta_{i,j}^{(k)}(x) \cdot n(x) > 0, \quad \forall x \in \partial \mathcal{O}_{i,j}^{(k)}, \quad k = 1, 2,$$

$$\zeta_{i,j}^{(k)}(x) \neq 0, \quad \forall x \in \Omega_{i,j} \setminus \mathcal{O}_{i,j}^{(k)}, \quad i = 1, 2.$$

Let $x_{i,j}^{(k)}$ be the integral curves of $\zeta_{i,j}^{(k)}$, i.e.

$$\begin{cases} \frac{dx_{i,j}^{(k)}(t)}{dt} = \zeta_{i,j}^{(k)}(x_{i,j}^{(k)}(t)), & t > 0, \\ x_{i,j}^{(k)}(0) = x_0, & x_0 \in S_{i,j}. \end{cases}$$

We also assume that there exists $\mathcal{T}_{i,j} > 0$ such that for all $x_0 \in S_{i,j}$, there exists $t_{i,j}^{(k)}(x_0) < \mathcal{T}_{i,j}$ satisfying

$$x_{i,j}^{(k)}(t) \in \Omega_{i,j} \setminus \mathcal{O}_{i,j}^{(k)}, \quad \text{for } 0 < t < t_{i,j}^{(k)}(x_0), \quad x_0 \in S_{i,j}, \quad k = 1, 2,$$

$$x_{i,j}^{(k)}(t_{i,j}^{(k)}(x_0)) \in \partial \mathcal{O}_{i,j}^{(k)}, \quad \text{for } x_0 \in S_{i,j}, \quad k = 1, 2.$$

In **case 2**, with an interface with say $\Omega_{i-1,k}$, we further assume that there exists a vector field $\zeta_{i,j} \in \mathcal{C}^1(\bar{\Omega}_{i,j}, \mathbb{R}^2)$, such that

$$\zeta_{i,j}(x) \cdot n(x) > 0, \quad \forall x \in S_{i,j}$$

$$\zeta_{i,j}(x) \cdot n(x) > 0, \quad \forall x \in S_{i-1,k}^{i,j}$$

$$\zeta_{i,j}(x) \neq 0, \quad \forall x \in \Omega_{i,j}$$

Let $x_{i,j}$ be the integral curves of $\zeta_{i,j}$, i.e.

$$\begin{cases} \frac{dx_{i,j}(t)}{dt} = \zeta_{i,j}(x_{i,j}(t)), & t > 0, \\ x_{i,j}(0) = x_0, & x_0 \in S_{i,j}. \end{cases}$$

We also assume that there exists $\mathcal{T}_{i,j} > 0$ such that for all $x_0 \in S_{i,j}$, there exists $t_{i,j}(x_0) < \mathcal{T}_{i,j}$ satisfying

$$x_{i,j}(t) \in \Omega_{i,j}, \quad \text{for } 0 < t < t_{i,j}(x_0), \quad x_0 \in S_{i,j},$$

$$x_{i,j}(t_{i,j}(x_0)) \in S_{i-1,k}^{i,j}, \quad \text{for } x_0 \in S_{i,j}.$$

Remark 5.2. In the simple case illustrated in Figure 3 where each open set Ω_i is connected, then only Ω_0 falls in the case 1 and $\Omega_1, \dots, \Omega_m$ satisfy case 2.

We let $A_1 = \cup_{i,j} \Omega_{i,j}$ where the union is performed over the i, j such that the sets $\Omega_{i,j}$ satisfy case 1 and $A_2 = \cup_{i,j} \Omega_{i,j}$ where the union is performed over the i, j such that the sets $\Omega_{i,j}$ satisfy case 2. Figures 3–6 illustrate cases where Assumption 5.1 is, or is not, satisfied.

We now assume

Assumption 5.3. The diffusion coefficient satisfies $c_i := c|_{\Omega_i} \in \mathcal{C}^1(\bar{\Omega}_i)$, $i = 0, \dots, m$.

Assumption 5.4. We have $0 < c_{\min} \leq c(x) \leq c_{\max}$, $x \in \Omega'$. If $R = S_{i,j}^{i+1,k} \neq \emptyset$, for $i = 0, \dots, m-1$, $1 \leq j \leq p_i$, and $1 \leq k \leq p_{i+1}$, then $c_i|_R \geq c_{i+1}|_R$.

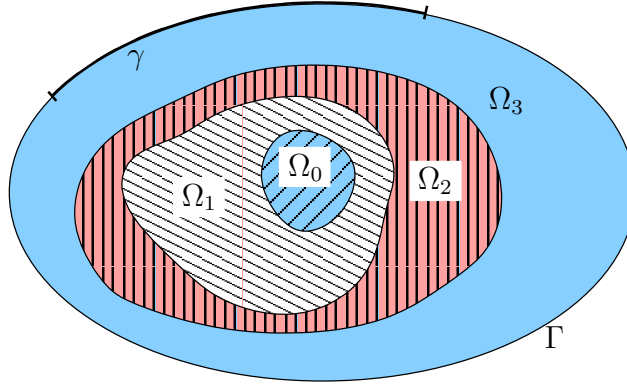


Figure 3: A geometric situation in which Assumption 5.1 is satisfied in the case $m = 3$.

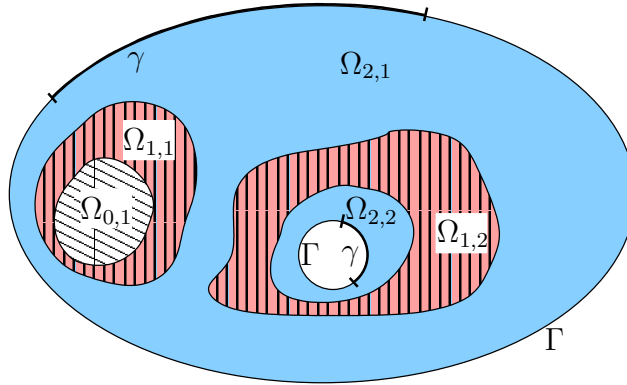


Figure 4: A geometric situation in which Assumption 5.1 is satisfied in the case $m = 2$.

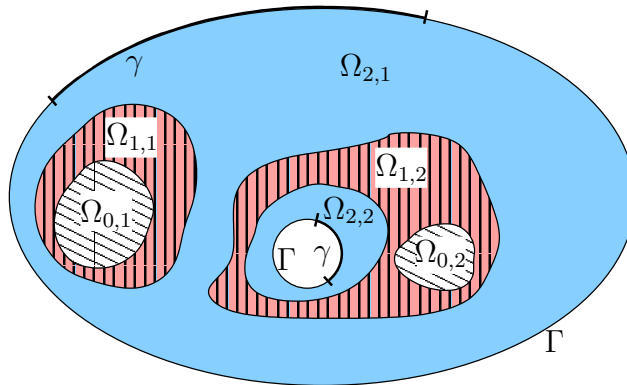


Figure 5: A geometric situation in which Assumption 5.1 is *not* satisfied.

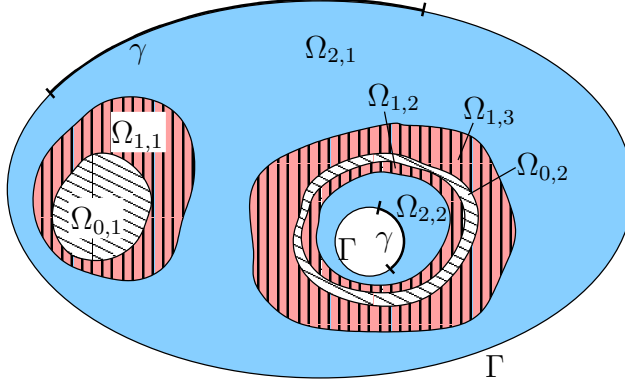


Figure 6: A geometric situation in which Assumption 5.1 is satisfied.

The geometric condition (GC_m) thus assumes that different materials are embedded in each other and the previous assumption thus states that the diffusion coefficient increases at the interfaces when one goes from the outer surface Γ to the inner parts of the material.

We can construct weight functions for a Carleman estimate with the following lemma.

Lemma 5.5. *Let the Geometric Condition (GC_m) be satisfied and let γ be a subset of $\Gamma = \partial\Omega$ satisfying Assumption 0.4. For the components $\Omega_{i,j}$ of Ω_i , $i = 1, \dots, m$, satisfying case 1 in (GC_m) we let $B_{i,j}^{(k)}$ and $\tilde{B}_{i,j}^{(k)}$, $k = 1, 2$ be open balls such that $B_{i,j}^{(k)} \subseteq \tilde{B}_{i,j}^{(k)} \subseteq \mathcal{O}_{i,j}^{(k)}$, $k = 1, 2$. Then, there exists two functions $\tilde{\beta}^{(1)}$ and $\tilde{\beta}^{(2)}$ in $\mathcal{C}^0(\bar{\Omega})$ such that*

$$\tilde{\beta}^{(k)}(x) = \begin{cases} \tilde{\beta}_1^{(k)} & \text{in } A_1, \\ \tilde{\beta}_2 & \text{in } \bar{A}_2, \end{cases}$$

and the functions $\tilde{\beta}_1^{(1)}$, $\tilde{\beta}_1^{(2)}$, and $\tilde{\beta}_2$ satisfy the following properties:

$$\tilde{\beta}_2 = 0 \text{ on } \Gamma \setminus \gamma, \quad \partial_n \tilde{\beta}_2 < 0 \text{ on } \Gamma \setminus \gamma.$$

For $\bar{\Omega}_{i,j}$ satisfying case 2, $\tilde{\beta}_2|_{\Omega_{i,j}} \in \mathcal{C}^2(\bar{\Omega}_{i,j})$, $\tilde{\beta}_2 > 0$ in $\Omega_{i,j}$, and

$$\tilde{\beta}_2 = i + 2 \text{ on } S_{i,j}, \quad \partial_n \tilde{\beta}_2 < 0 \text{ on } S_{i,j}, \text{ if } i < m,$$

and

$$|\nabla \tilde{\beta}_2| > 0 \text{ on } \bar{\Omega}_{i,j}.$$

For $\bar{\Omega}_{i,j}$ satisfying case 1, for $k = 1, 2$, $\tilde{\beta}_1^{(k)}|_{\Omega_{i,j}} \in \mathcal{C}^2(\bar{\Omega}_{i,j})$, $\tilde{\beta}_1^{(k)} > 0$ in $\Omega_{i,j}$,

$$\tilde{\beta}_1^{(k)} = 2 \text{ on } S_{i,j}, \quad k = 1, 2,$$

$$\tilde{\beta}_1^{(1)} \geq 2\tilde{\beta}_1^{(2)} \text{ in } \tilde{B}_{i,j}^{(2)},$$

$$\tilde{\beta}_1^{(2)} \geq 2\tilde{\beta}_1^{(1)} \text{ in } \tilde{B}_{i,j}^{(1)},$$

and

$$|\nabla \tilde{\beta}_1^{(k)}| > 0 \text{ in } \bar{\Omega}_{i,j} \setminus B_{i,j}^{(k)}, \quad k = 1, 2.$$

Finally on an (non-empty) interface $R = S_{i,j}^{i+1,k}$

$$(c_i \partial_n \tilde{\beta}^{(k)})|_{\Omega_{i,j}} = (c_{i+1} \partial_n \tilde{\beta}^{(k)})|_{\Omega_{i+1,k}} \text{ on } R, \quad k = 1, 2.$$

We set $\mathcal{O}^{(k)} = \bigcup_{i,j} \mathcal{O}_{i,j}^{(k)}$, $k = 1, 2$, where the union is performed over the i, j such that the sets $\Omega_{i,j}$ satisfy case 1 in Assumption 5.1.

We introduce the functions $\beta^{(k)}$ and the weight functions $\varphi^{(k)}, \eta^{(k)}$, $k = 1, 2$, following section 1. Let $g \in H^1([t_0, T], H^{\frac{1}{2}}(S))$. We introduce transmission condition (TC3) on the interval $[t_0, T]$: for each (non-empty) interface $R = S_{i,j}^{i+1,k}$, we have

$$\begin{aligned} \text{(TC3)} \quad & q|_{\Omega_{i,j} \times [t_0, T]} = q|_{\Omega_{i+1,k} \times [t_0, T]}, \text{ on } R \\ & (c_i \partial_n q)|_{\Omega_{i,j} \times [t_0, T]} = (c_{i+1} \partial_n q)|_{\Omega_{i+1,k} \times [t_0, T]} + g(x, t), \text{ on } R \end{aligned}$$

for a function q which is H^2 in each open set $\Omega_{i,j}$, $i = 0, \dots, m$, $j = 1, \dots, p_i$. Following Section 1 we introduce

$$\begin{aligned} \mathfrak{N}_g = \Big\{ & q \in H^1(t_0, T, H_0^1(\Omega)); \quad q|_{(t_0, T) \times \Omega_i} \in L^2(t_0, T, H^2(\Omega_i)), \quad i = 0, 1, \\ & q|_{\Sigma} = 0 \text{ and } q \text{ satisfies (TC3) a.e. w.r.t. } t \Big\}. \end{aligned}$$

We then obtain

Theorem 5.6. *Let γ be a subset of the boundary Γ of an open set Ω of \mathbb{R}^n that satisfies Condition (GC_m) , and γ satisfies Assumption 0.4. Assume further that there exists $\Delta > 0$ such that, for each (non-empty) interface $R = S_{i,j}^{i+1,k}$, we have $c_i|_R - c_{i+1}|_R \geq \Delta > 0$. Let $g \in H^1(t_0, T, H^{\frac{1}{2}}(S))$. There exists $\lambda_1 = \lambda_1(\Omega, \gamma, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, c_{\min}, c_{\max}, \Delta) > 0$, $s_1 = s_1(\lambda_1) > 0$ and a positive constant $C = C(\Omega, \gamma, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, c_{\min}, c_{\max}, \Delta)$ so that Carleman estimate (1.5) holds for $s \geq s_1$, $\lambda \geq \lambda_1$ and for all $q \in \mathfrak{N}_g$.*

Following the method of Section 2 we thus obtain the following stability result.

Theorem 5.7. *Let γ be a subset of the boundary Γ of an open set Ω of \mathbb{R}^n that satisfies Condition (GC_m) , and γ satisfies Assumption 0.4. We assume that the diffusion coefficients c and \tilde{c} satisfy Assumptions 5.4 and that c and \tilde{c} are piecewise constant, in the sense that $c|_{\Omega_i}$, resp. $\tilde{c}|_{\Omega_i}$, are constant in each connected component of Ω_i , $i = 0, \dots, m$ and that furthermore there exists $\Delta > 0$ such that, for each (non-empty) interface $R = S_{i,j}^{i+1,k}$, we have $c_i|_R - c_{i+1}|_R \geq \Delta > 0$. Let y_0, \tilde{y}_0 in $L^2(\Omega)$ and let y, \tilde{y} be solutions to (2.1)–(2.2) satisfying Assumption 2.1. Then there exists a constant C*

$$C = C(\Omega, T, t_0, \gamma, S, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, M, K, r, c_{\min}, c_{\max}, \Delta)$$

such that

$$|c - \tilde{c}|_{L^\infty(\Omega)}^2 \leq C |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0, T) \times \gamma)}^2 + C |\Delta y(T', \cdot) - \Delta \tilde{y}(T', \cdot)|_{L^2(\Omega')}^2.$$

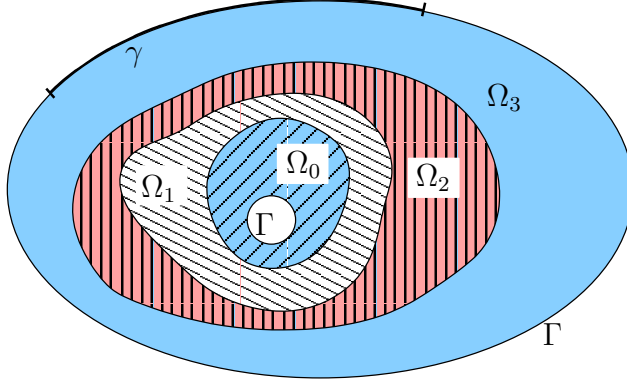


Figure 7: A geometrical situation for which only one set of weight functions is needed to obtain a Carleman estimate.

Similarly, following Section 4 we obtain

Theorem 5.8. *Under the hypothesis of Theorem 5.7 in addition to Assumption 4.1 there exists a constant $C > 0$*

$$C = C(\Omega, T, t_0, \gamma, S, \mathcal{O}^{(1)}, \mathcal{O}^{(2)}, M, K, r, c_{\min}, c_{\max}, \Delta)$$

such that

$$|y_0 - \tilde{y}_0|_{L^2(\Omega)} \leq C / \ln \left(|(y - \tilde{y})(T')|_{H^2(\Omega')} + |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 \right)$$

for $|(y - \tilde{y})(T')|_{H^2(\Omega')} + |\partial_n(\partial_t y - \partial_t \tilde{y})|_{L^2((0,T) \times \gamma)}^2 < 1$.

Remark 5.9. In the generalization proposed here we assume some monotonicity for the diffusion coefficient c at the inner interfaces. Note that if some of these assumptions are not satisfied we may still obtain stability results by introducing additional measurements in some $\omega \times (t_0, T)$ for $\omega \subset \Omega$. Each geometrical configuration requires the derivation of a particular Carleman estimate to apply the method exposed here.

Remark 5.10. Here we have assumed

$$V_0 \Subset V_1 \Subset \dots \Subset V_m = \Omega.$$

Observe that this is not true in the case illustrated in Figure 7. then ∂_Ω contains a component of Γ . In such a case we can however obtain a Carleman estimate by using a single set of weight functions η and φ instead of two (see (1.4) and (1.5)). Stability results follow similarly in such a geometrical configuration.

A Basic regularity properties

Let \mathcal{A} be formally defined by $-\nabla \cdot (c \nabla (\cdot))$ on $L^2(\Omega)$. The diffusion coefficient c is first assumed to be in $L^\infty(\Omega)$ and such that $c(x) \geq \alpha > 0$ for all $x \in \Omega$. We

denote by A the unbounded operator with domain

$$\mathcal{D}_A = \{u \in H_0^1(\Omega); \nabla \cdot (c \nabla u) \in L^2(\Omega)\}.$$

defined by $A(u) = -\nabla \cdot (c \nabla(u))$ for $u \in \mathcal{D}_A$.

Lemma A.1. *The operator A is maximal accretive and self-adjoint on $L^2(\Omega)$.*

Proof. Let $u \in \mathcal{D}_A$. We have

$$(Au, u)_{L^2(\Omega)} = - \int_{\Omega} \nabla \cdot (c \nabla u) u \, dx = \int_{\Omega} c \nabla u \cdot \nabla u \, dx \geq 0$$

since $u \in H_0^1(\Omega)$. Hence A is accretive. To prove that A is maximal accretive on $L^2(\Omega)$ we have to prove that $R(I + A) = L^2(\Omega)$. Let $f \in L^2(\Omega)$. We consider the elliptic problem

$$\begin{cases} -\nabla \cdot (c \nabla u) + u = f, \\ u|_{\partial\Omega} = 0. \end{cases}$$

We seek a weak solution, i.e. a function u satisfying,

$$(A.1) \quad \int_{\Omega} c \nabla u \cdot \nabla \varphi \, dx + \int_{\Omega} u \varphi \, dx = \int_{\Omega} f \varphi \, dx, \quad \forall \varphi \in H_0^1(\Omega).$$

We define the bilinear form $a(u, v) = \int_{\Omega} c \nabla u \cdot \nabla v \, dx + \int_{\Omega} uv \, dx$. Observe that a is continuous and coercive since $c \geq \alpha > 0$. The Lax-Milgram theorem thus yields a unique solution u to (A.1) which is in $H_0^1(\Omega)$. Observe from (A.1) that for all $\varphi \in H_0^1(\Omega)$ we have

$$|\int_{\Omega} c \nabla u \cdot \nabla \varphi \, dx| \leq C |\varphi|_{L^2(\Omega)}.$$

This exactly means that $\nabla \cdot (c \nabla u) \in L^2(\Omega)$. The solution u to (A.1) is thus in \mathcal{D}_A . Therefore, we obtain $Au + u = f$ in $L^2(\Omega)$ (use Corollary 2.1 in [11]). Hence $R(I + A) = L^2(\Omega)$.

Now, A is symmetric which implies that it is self-adjoint [5, Proposition VII.6]. ■

Proposition A.2. *Let $u_0 \in L^2(\Omega)$. There exists a unique u such that*

$$u \in \mathcal{C}([0, T]; L^2(\Omega)) \cap \mathcal{C}^1([0, T]; L^2(\Omega)) \cap \mathcal{C}([0, T]; \mathcal{D}_A),$$

and

$$(A.2) \quad \begin{cases} \partial_t u - \nabla \cdot (c \nabla(u)) = 0, & t \in]0, T], \\ u(0) = u_0. \end{cases}$$

for $T > 0$. (T can be chosen to be ∞). Furthermore $u \in L^2(0, T; H_0^1(\Omega))$ and $u \in \mathcal{C}^k([0, T]; \mathcal{D}_{A^l})$ for all $k, l \in \mathbb{N}$.

If $u_0 \in \mathcal{D}_A$ then

$$u \in \mathcal{C}^1([0, T]; L^2(\Omega)) \cap \mathcal{C}([0, T]; \mathcal{D}_A).$$

Proof. With Lemma A.1, we apply the Hille-Yoshida theorem for a self-adjoint operator on a Hilbert space (see e.g. [5]). ■

In the case $u_0 \in \mathcal{D}_A$ then $\partial_t u(0) = -Au_0 \in L^2(\Omega)$ [21]. $\partial_t u$ is thus solution to (A.2) with $-Au_0$ for initial condition.

Proposition A.3. *The semigroup $S(t)$ generated by the unbounded operator A on $L^2(\Omega)$ is analytic.*

Proof. From Theorem VII.7 in [5] and Lemma A.1 above, we obtain that $A \circ S(t)$ is bounded if $t > 0$ and $\|A \circ S(t)\| \leq \frac{1}{t}$. Since 0 is in the resolvent set of A , we can apply Theorem 2.5.2 in [21] which yields the result. ■

We have obtained so far that if u is the solution to (A.2) with an initial condition in $L^2(\Omega)$ then, $u(t, \cdot) \in H_0^1(\Omega)$ and $\nabla \cdot (c \nabla(u)) \in L^2(\Omega)$, if $t > 0$. If the diffusion coefficient c is piecewise \mathcal{C}^1 and discontinuous across some \mathcal{C}^2 interface S we can then assert that $c \partial_n u|_S$ is well defined as an element of $H^{-\frac{1}{2}}(S)$ [11, Theorem 2.2], where $\partial_n u|_S = (\nabla u)|_S \cdot n$ with n denoting the normal unit vector to S .

We now give further regularity properties when placed in the geometrical situation studied in this paper. Let $\Omega \subset \mathbb{R}^n$ be a bounded connected open set. The set $\overline{\Omega}$ is assumed to be a \mathcal{C}^2 submanifold with boundary in \mathbb{R}^n . Let Ω_0 and Ω_1 be two non-empty open subsets of Ω such that

$$\Omega_0 \Subset \Omega, \text{ and } \Omega_1 = \Omega \setminus \overline{\Omega}_0.$$

We denote $S = \overline{\Omega}_0 \cap \overline{\Omega}_1$ the interface, which is assumed to be \mathcal{C}^2 and we denote n the outward unit normal to Ω_1 on S and also the outward unit normal to Ω on Γ . Let S_0 (resp. S_1) be the side of the interface S corresponding to the positive (resp. negative) direction of the normal n .

Proposition A.4. *Let the diffusion coefficient, c , be such that $c|_{\Omega_i} \in \mathcal{C}^1(\overline{\Omega}_i)$, $i = 0, 1$. If $p \in \mathcal{D}_A$ then $p|_{\Omega_i} \in H^2(\Omega_i)$, $i = 0, 1$. Furthermore $|p|_{\Omega_i}|_{H^2(\Omega_i)} \leq C|\nabla \cdot (c \nabla p)|_{L^2(\Omega)}$.*

We introduce some notations to be used in the proof of the proposition. Let us write $y \in \mathbb{R}^n$ in the form $y = (y', y_n)$ with $y' \in \mathbb{R}^{n-1}$ and y_n and define

$$\begin{aligned} \mathbb{R}_+^n &= \{y = (y', y_n); y_n > 0\}, \\ \mathbb{R}_-^n &= \{y = (y', y_n); y_n < 0\}, \\ V &= \{y = (y', y_n); |y'| < 1 \text{ and } |y_n| < 1\}, \\ V_+ &= V \cap \mathbb{R}_+^n, \\ V_- &= V \cap \mathbb{R}_-^n, \\ V_0 &= \{y = (y', 0); y \in V\}. \end{aligned}$$

Proof. The proof is along the line of that of Theorem IX.25 in [5] and makes use of Nirenberg's translation method. We set $f := -\nabla \cdot (c \nabla p) \in L^2(\Omega)$. We cover Ω with a finite number of sufficiently small open balls, \mathcal{B}_j , $j = 1, \dots, N$ such that either 1) $\mathcal{B}_j \Subset \Omega_0$, 2) $\mathcal{B}_j \Subset \Omega_1$, 3) $\Omega_0 \cap \mathcal{B}_j = \emptyset$ and part of $\Gamma = \partial\Omega$ is in \mathcal{B}_j , or 4) $\Omega_i \cap \mathcal{B}_j \neq \emptyset$, $i = 0, 1$, and part of S is in \mathcal{B}_j . Let θ_j , $j = 1, \dots, N$ be a \mathcal{C}^2 partition of unity subordinated to the covering \mathcal{B}_j , $j = 1, \dots, N$.

Lemma A.5. *The open balls \mathcal{B}_i and the functions θ_i can be chosen such that $\partial_n \theta_i|_S = 0$, $i = 1, \dots, N$*

Proof. In a small neighborhood of S we can extend the unit normal vector n to S into a \mathcal{C}^2 vector field v . In an even smaller neighborhood U of S we can assume that v is such that $|v| \geq a > 0$. If we integrate this vector field we find that there is $\epsilon > 0$ such that the flow χ_σ of this \mathcal{C}^2 vector field over the interval $]-\epsilon, \epsilon[$ is confined in U , since S is compact. For $y \in S$, the orientation of the unit normal (see above) is such that $\chi_\sigma(y) \in \Omega_1$ for $\sigma \in]-\epsilon, 0[$, $\chi_0(y) = y$ and $\chi_\sigma(y) \in \Omega_0$ for $\sigma \in]0, \epsilon[$. Set now

$$W = \{\chi_\sigma(y); y \in S \text{ and } \sigma \in]-\epsilon, \epsilon[\}$$

which is an open neighborhood of S . Note that if x in W then there exists a unique $y \in S$ and a unique $\sigma \in]-\epsilon, \epsilon[$ such that $x = \chi_\sigma(y)$.

Define $\delta = \text{dist}(S, \Omega \setminus W)$ which is positive. Cover $\overline{\Omega \setminus W}$ with a finite number of open balls of diameter $\delta/2$, say $\mathcal{B}_1, \dots, \mathcal{B}_k$. Thus $\mathcal{B}_1, \dots, \mathcal{B}_k, W$ is an open covering of $\overline{\Omega}$. Note that there exists a neighborhood of S , contained in W , which does not intersect the open balls $\mathcal{B}_1, \dots, \mathcal{B}_k$. Let $\theta_1, \dots, \theta_k, \theta$ be a partition of unity subordinated to this covering:

$$\begin{aligned} \theta_i &\in \mathcal{C}_c^\infty(\mathcal{B}_i), \quad i = 1, \dots, k, \quad \theta \in \mathcal{C}_c^\infty(W), \\ (\theta + \sum_{1 \leq i \leq k} \theta_i)|_{\overline{\Omega}} &= 1. \end{aligned}$$

Observe that θ is equal to 1 in a neighborhood of S .

Let \mathcal{C}_j , $j = 1, \dots, r$, be open balls that form an open covering of S in Ω and set $\tilde{\mathcal{C}}_j = S \cap \mathcal{C}_j$, $j = 1, \dots, r$, which is an open covering of S (for the topology induced on S). Let $\tilde{\psi}_j$, $j = 1, \dots, r$, be a partition of unity subordinated to this latter covering and define now

$$\begin{aligned} \psi_j(x) &= \tilde{\psi}_j(y)\theta(x), \quad x \in W, \quad j = 1, \dots, r, \\ &\text{where } x = \chi_\sigma(y), \text{ with } y \in S, \quad \sigma \in]-\epsilon, \epsilon[. \end{aligned}$$

Note that $\sum \psi_j = \theta$. Furthermore, $\partial_n \psi_j|_S = \tilde{\psi}_j \partial_n \theta|_S = 0$. Let $N = k + r$. The functions $\theta_1, \dots, \theta_k, \psi_1, \dots, \psi_r$ form a partition of unity that satisfies the required property. \blacksquare

Continuation of the proof of Proposition A.4. We shall prove that for all $j = 1, \dots, N$ the function $\theta_j p|_{\Omega_i}$ is in $H^2(\Omega_i)$, $i = 0, 1$. This yields the result. The cases 1, 2, and 3 described above follow from the proof of Theorem IX.25 in [5]. We thus focus our attention to case 4, i.e., when the considered ball \mathcal{B}_i contains part of the interface S . For simplicity we shall write θ, \mathcal{B} instead of θ_i and \mathcal{B}_i .

Let $v = \theta p$. Observe that $c \nabla v = c(\nabla \theta)p + c\theta(\nabla p)$ and that, for $\psi \in H_0^1(\mathcal{B})$,

$$\begin{aligned} \int_{\mathcal{B}} c \nabla v \nabla \psi \, dx &= \int_{\mathcal{B}} c(\nabla \theta)p \nabla \psi \, dx + \int_{\mathcal{B}} c(\nabla p)\theta \nabla \psi \, dx \\ &= - \sum_{i=0,1} \int_{\mathcal{B} \cap \Omega_i} \nabla(c(\nabla \theta)p) \psi \, dx + \int_{\mathcal{B}} c \nabla p \nabla(\theta \psi) \, dx - \int_{\mathcal{B}} c(\nabla p \nabla \theta) \psi \, dx, \end{aligned}$$

since $\theta \nabla \psi = \nabla(\theta \psi) - (\nabla \theta) \psi$. The integration by parts is justified since $c \nabla \theta|_{\Omega_i} \in \mathcal{C}^1(\overline{\Omega_i \cap \mathcal{B}})$, $i = 0, 1$, $p \in H^1(\mathcal{B})$ and $c \partial_n \theta|_S = 0$. This finally yields

$$\int_{\mathcal{B}} c \nabla v \nabla \psi \, dx = \int_{\mathcal{B}} g \psi \, dx,$$

where $g = -\nabla(c \nabla \theta) p + \theta f - 2c(\nabla p \nabla \theta)$ on Ω_i , $i = 0, 1$. We have $g \in L^2(\mathcal{B})$ since $(c \nabla \theta)|_{\Omega_i}$ is in $\mathcal{C}^1(\overline{\Omega_i \cap \mathcal{B}})$, $i = 0, 1$. Observe that $|g|_{L^2(\mathcal{B})} \leq C|f|_{L^2(\Omega)}$ since $|p|_{H_0^1(\Omega)} \leq C|f|_{L^2(\Omega)}$.

Because of the regularity imposed on the interface S , there exists a \mathcal{C}^2 diffeomorphism $J : \mathcal{B} \rightarrow V$ such that $J(\Omega_0 \cap \mathcal{B}) = V_-$, $J(\Omega_1 \cap \mathcal{B}) = V_+$, and $J(S \cap \mathcal{B}) = V_0$. Let $H = J^{-1}$. (In fact, note that in the proof of Lemma A.5 the neighborhoods \mathcal{C}_j , $j = 1, \dots, r$ and $\epsilon > 0$ can be chosen sufficiently small so that such a diffeomorphism exists.)

Let $w = v \circ H$. Then we can prove that w satisfies a problem of the form

$$(A.3) \quad \sum_{k,l} \int_V \underline{c}_{kl} \partial_{x_k} w \partial_{x_l} \psi \, dx = \int_V \underline{g} \psi \, dx, \quad \psi \in H_0^1(V),$$

where \underline{c} satisfies the uniform ellipticity condition $\sum_{k,l} \underline{c}_{kl} \xi_k \xi_l \geq \underline{\alpha} |\xi|^2$ with $\underline{\alpha} > 0$, and $\underline{g} \in L^2(V)$ (adapt Lemma IX.8 in [5]). The functions \underline{c}_{kl} are piecewise \mathcal{C}^1 with a discontinuity across the interface V_0 . Note that $w \in H_0^1(V)$ with its support finitely away from ∂V .

Let h be parallel to V_0 . Define D_h by $D_h(\rho) = (\rho(x+h) - \rho(x))/|h|$. Observe that $D_{-h}(D_h w) \in H_0^1(V)$ for $|h|$ sufficiently small and set $\psi = D_{-h}(D_h w)$. This yields

$$(A.4) \quad \sum_{k,l} \int_V D_h(\underline{c}_{kl} \partial_{x_k} w) \partial_{x_l} (D_h w) \, dx = \sum_{k,l} \int_V [\underline{c}_{kl}(y+h) D_h(\partial_{x_k} w) \\ + (D_h \underline{c}_{kl}) \partial_{x_k} w] \partial_{x_l} (D_h w) \, dx = \int_V \underline{g} D_{-h}(D_h w) \, dx.$$

We Note that

$$\left| \int_V (D_h \underline{c}_{kl}) \partial_{x_k} w \partial_{x_l} (D_h w) \, dx \right| \leq C|w|_{H^1(V)} |\nabla(D_h w)|_{L^2(V)}, \quad 1 \leq k, l \leq n,$$

and

$$\int_V \underline{g} D_{-h}(D_h w) \, dx \leq |\underline{g}|_{L^2(V)} |D_{-h}(D_h w)|_{L^2(V)} \leq |\underline{g}|_{L^2(V)} |\nabla(D_h w)|_{L^2(V)}$$

(Recall that if $\rho \in H_0^1(V)$ with its support finitely away from the boundary ∂V then $|D_h(\rho)|_{L^2(V)} \leq |\nabla \rho|_{L^2(V)}$ for $|h|$ sufficiently small [5, Proposition IX.3].)

From (A.4) we find

$$\underline{\alpha} |\nabla(D_h w)|_{L^2(V)}^2 \leq C(|w|_{H^1(V)} + |\underline{g}|_{L^2(V)}) |\nabla(D_h w)|_{L^2(V)}.$$

With $\psi = w$ we find $|\nabla w|_{L^2(V)} \leq C|\underline{g}|_{L^2(V)}$ which yields $|\nabla(D_h w)|_{L^2(V)} \leq C|\underline{g}|_{L^2(V)}$. Thus for $\psi \in H_0^1(V)$,

$$\begin{aligned} \left| \int_V (\partial_{x_l} w)(D_{-h} \psi) dx \right| &= \left| \int_V (D_h w)(\partial_{x_l} \psi) dx \right| = \left| \int_V \partial_{x_l} (D_h w) \psi dx \right| \\ &\leq C|\underline{g}|_{L^2(V)} |\psi|_{L^2(V)}. \end{aligned}$$

Letting $|h|$ go to zero we obtain

$$\left| \int_V \partial_{x_k} w \partial_{x_l} \psi dx \right| \leq C|\underline{g}|_{L^2(V)} |\psi|_{L^2(V)}, \quad (k, l) \neq (n, n),$$

for all $\psi \in H_0^1(V)$. Hence $\partial_{x_k} \partial_{x_l} w \in L^2(V)$ if $(k, l) \neq (n, n)$. Let us now choose $\psi \in H_0^1(V_+)$ in (A.3). We obtain that in V_+

$$\partial_{x_n}^2 w = -\frac{1}{\underline{c}_{nn}} \left(f + \partial_{x_k} (\underline{c}_{kl} \partial_{x_l} w) + (\partial_{x_n} \underline{c}_{nn}) \partial_{x_n} w \right),$$

since $\underline{c}_{nn} \geq \alpha > 0$, which yields $\partial_{x_n}^2 w|_{V_+} \in L^2(V_+)$. Finally $w|_{V_+} \in H^2(V_+)$ and $|w|_{V_+}|_{H^2(V_+)} \leq C|\underline{g}|_{L^2(V)}$. We proceed in a similar fashion to prove $w|_{V_-} \in H^2(V_-)$ and $|w|_{V_-}|_{H^2(V_-)} \leq C|\underline{g}|_{L^2(V)}$. \blacksquare

Corollary A.6. *Let the diffusion coefficient, c , be such that $c|_{\Omega_i} \in \mathcal{C}^1(\overline{\Omega_i})$, $i = 0, 1$. If $u_0 \in L^2(\Omega)$ then the solution u to (A.2) is such that $u|_{\Omega_i}(t, \cdot) \in H^2(\Omega_i)$, $i = 0, 1$, when $t > 0$. Furthermore $|u|_{\Omega_i}(t, \cdot)|_{H^2(\Omega_i)} \leq C|\partial_t u(t, \cdot)|_{L^2(\Omega)}$.*

Proof. Observe that $u \in \mathcal{D}_A$ by Proposition A.2 if $t > 0$. Hence $f := -\partial_t u = -\nabla(c\nabla u) \in L^2(\Omega)$ if $t > 0$. \blacksquare

Remark A.7. Observe that with the previous proposition the normal trace of $c\nabla u$ on S , i.e. $c\partial_n u|_S$, is now found to belong to $H^{\frac{1}{2}}(S)$.

Requiring additionnal smoothness for the diffusion coefficient, c , in the open sets Ω_i , $i = 0, 1$, and for the boundary $\partial\Omega$ and the interface S we can improve the previous result.

Proposition A.8. *Let $m \in \mathbb{N}$ and let $c|_{\Omega_i} \in \mathcal{C}^{m+1}(\overline{\Omega_i})$, $i = 0, 1$, and let S and $\partial\Omega$ be of class \mathcal{C}^{m+2} . If $p \in \mathcal{D}_A$ and $A(p) = f$ with $f|_{\Omega_i} \in H^m(\Omega_i)$ then $p|_{\Omega_i} \in H^{m+2}(\Omega_i)$, $i = 0, 1$. Furthermore $|p|_{\Omega_i}|_{H^{m+2}(\Omega_i)} \leq C|f|_{\Omega_i}|_{H^m(\Omega_i)}$, $i = 0, 1$.*

Proof. We proceed by induction. The case $m = 0$ is contained in the proof of Proposition A.4. Let $m_0 > 0$. Assume the result is true for $0 \leq m \leq m_0 - 1$ and assume $f|_{\Omega_i} \in H^{m_0}(\Omega_i)$, $i = 0, 1$. Thus, $f|_{\Omega_i} \in H^{m_0-1}(\Omega_i)$ which yields $p|_{\Omega_i} \in H^{m_0+1}(\Omega_i)$ by induction, and $|p|_{\Omega_i}|_{H^{m_0+1}(\Omega_i)} \leq C|f|_{\Omega_i}|_{H^{m_0-1}(\Omega_i)} \leq C|f|_{\Omega_i}|_{H^{m_0}(\Omega_i)}$, $i = 0, 1$.

We use the same partition of unity θ_j , $j = 1, \dots, N$, as in the proof of Proposition A.4 and prove the result for $v := \theta_i p$, i.e. $v|_{\Omega_i} \in H^{m_0+2}(\Omega_i)$, $i = 0, 1$. We choose here the functions θ_i to be \mathcal{C}^{2m_0+2} . We focus our attention to the

case where $\text{supp}(\theta_i)$ contains part of S since the other cases follow from classical results. We obtain

$$\int_{\mathcal{B}} c \nabla v \nabla \psi \, dx = \int_{\mathcal{B}} g \psi \, dx, \quad \psi \in H_0^1(\mathcal{B}),$$

where $g = -\nabla(c \nabla \theta) p + \theta f - 2c(\nabla p \nabla \theta)$ on Ω_i , $i = 0, 1$. Observe that $g|_{\Omega_i} \in H^{m_0}(\Omega_i)$ and $|g|_{\Omega_i}|_{H^{m_0}(\Omega_i)} \leq |f|_{\Omega_i}|_{H^{m_0}(\Omega_i)}$, $i = 0, 1$.

With the notations of the proof of Proposition A.4 let $w = v \circ H$. Then w satisfies

$$(A.5) \quad \sum_{k,l} \int_V \underline{c}_{kl} \partial_{x_k} w \partial_{x_l} \psi \, dx = \int_V \underline{g} \psi \, dx, \quad \psi \in H_0^1(V),$$

where \underline{c} is uniformly elliptic and \mathcal{C}^{m+1} in each Ω_i , $i = 0, 1$, and $\underline{g}|_{V_{\pm}} \in H^{m_0}(V_{\pm})$ and

$$\begin{aligned} |\underline{g}|_{V_+}|_{H^{m_0}(V_+)} &\leq |f|_{\Omega_1}|_{H^{m_0}(\Omega_1)}, \\ |\underline{g}|_{V_-}|_{H^{m_0}(V_-)} &\leq |f|_{\Omega_0}|_{H^{m_0}(\Omega_0)}. \end{aligned}$$

As in the proof of Lemma IX.7 in [5] we can show $\partial_{x_j} w \in H_0^1(V)$ for $j = 1, \dots, n-1$. Choosing $\psi = \partial_{x_j} \varphi$ with $\varphi \in \mathcal{C}_c^\infty(V)$ we obtain

$$\sum_{k,l} \int_V \underline{c}_{kl} \partial_{x_k} (\partial_{x_j} w) \partial_{x_l} \varphi \, dx = \int_V \tilde{g} \varphi \, dx, \quad \varphi \in \mathcal{C}_c^\infty(V).$$

with $\tilde{g}|_{V_{\pm}} \in H^{m_0-1}(V_{\pm})$ satisfying $|\tilde{g}|_{V_{\pm}}|_{H^{m_0-1}(V_{\pm})} \leq C |\underline{g}|_{V_{\pm}}|_{H^{m_0}(V_{\pm})}$. From the induction we obtain $(\partial_{x_j} w)|_{V_{\pm}} \in H^{m_0+1}(V_{\pm})$, $j = 1, \dots, n-1$ with

$$|(\partial_{x_j} w)|_{V_{\pm}}|_{H^{m_0+1}(V_{\pm})} \leq |\underline{g}|_{V_{\pm}}|_{H^{m_0}(V_{\pm})}$$

(here, we implicitly proceed with an induction proof for the regularity of v). Since $w|_{V_{\pm}} \in H^{m_0+1}(V_{\pm})$, to conclude we have to prove

$$(A.6) \quad \left| \int_{V_{\pm}} \partial_{x_n} w \partial^\alpha \psi \, dx \right| \leq C |\underline{g}|_{V_{\pm}}|_{H^{m_0}(V_{\pm})} |\psi|_{L^2(V_{\pm})}, \quad \psi \in \mathcal{C}_c^\infty(V_{\pm}), \quad |\alpha| = m_0 + 1,$$

where α is a multi-index $\alpha = (\alpha_1, \dots, \alpha_n)$. We first treat the case $\alpha_j \neq 0$ for some j in $\{1, \dots, n-1\}$, say $\alpha_1 \neq 0$ then

$$\int_{V_{\pm}} \partial_{x_n} w \partial^\alpha \psi \, dx = \int_{V_{\pm}} \partial_{x_1} w \partial^{\alpha'} \psi \, dx, \quad \text{with } \alpha' = (\alpha_1 - 1, \alpha_2, \dots, \alpha_{n-1}, 1)$$

since $\psi \in \mathcal{C}_c^\infty(V_{\pm})$ and we obtain estimate (A.6) since $\partial_{x_1} w|_{V_{\pm}} \in H^{m_0+1}(V_{\pm})$ and $|\alpha'| = m_0 + 1$.

We write in V_{\pm}

$$\partial_{x_n}^2 w = -\frac{1}{\underline{c}_{nn}} \left(f + \partial_{x_k} (\underline{c}_{kl} \partial_{x_l} w) + (\partial_{x_n} \underline{c}_{nn}) \partial_{x_n} w \right),$$

and find $\partial_{x_n}^2 w|_{V_{\pm}} \in H^{m_0}(V_{\pm})$ which yields estimate (A.6) in the case $\alpha = (0, \dots, 0, m_0 + 1)$. \blacksquare

Corollary A.9. *Let $l \in \mathbb{N}$ and let $c_{|\Omega_i} \in \mathcal{C}^{2l+1}(\Omega_i)$, $i = 0, 1$, and S and $\partial\Omega$ be of class \mathcal{C}^{2l+2} . If $p \in \mathcal{D}_{A^{l+1}}$ then $p|_{\Omega_i} \in H^{2l+2}(\Omega_i)$ and $|p|_{\Omega_i}|_{H^{2l+2}(\Omega_i)}| \leq C|A^{l+1}(p)|_{L^2(\Omega)}$, $i = 0, 1$.*

Proof. We proceed by induction. The case $l = 0$ is contained in Proposition A.4. Assume the result is true for $l \leq l_0 - 1$ and assume $p \in \mathcal{D}_{A^{l_0+1}}$. Thus $f := A(p) \in \mathcal{D}_{A^{l_0}}$ and thus $f|_{\Omega_i} \in H^{2l_0}(\Omega_i)$, and $|f|_{\Omega_i}|_{H^{2l_0}(\Omega_i)}| \leq C|A^{l_0+1}(p)|_{L^2(\Omega)}$, $i = 0, 1$. Finally $p|_{\Omega_i} \in H^{2l_0+2}(\Omega_i)$ and $|p|_{\Omega_i}|_{H^{2l_0+2}(\Omega_i)}| \leq C|A^{l_0+1}(p)|_{L^2(\Omega)}$, $i = 0, 1$, by Proposition A.8. ■

Remark A.10. Observe that the constant in the inequalities of the previous results can be chosen uniform w.r.t. c for $c_{min} \leq c \leq c_{max}$.

Combining the results just obtained and those of Proposition A.2 we obtain

Corollary A.11. *Let $m \in \mathbb{N}$ and let $c_{|\Omega_i} \in \mathcal{C}^{m+1}(\Omega_i)$, $i = 0, 1$, and S and $\partial\Omega$ be of class \mathcal{C}^{m+2} . Then if $u_0 \in L^2(\Omega)$ the solution u to (A.2) is such that $u|_{\Omega_i} \in \mathcal{C}^k((0, T], H^{m+2}(\Omega_i))$, $i = 0, 1$, for all $k \in \mathbb{N}$.*

Proof. Let $k \in \mathbb{N}$ and let $l \in \mathbb{N}$ be such that $2l \leq m \leq 2l + 2$. Since $u \in \mathcal{C}^k((0, T], \mathcal{D}_{A^{l+1}})$ then $u|_{\Omega_i} \in \mathcal{C}^k((0, T], H^{2l+2}(\Omega_i))$ by Corollary A.9 since $c_{|\Omega_i} \in \mathcal{C}^{2l+1}(\Omega_i)$. In the case $m = 2l$ we have obtained the sought result.

In the case $m = 2l + 1$ we similarly find $Au|_{\Omega_i} \in \mathcal{C}^k((0, T], H^{m+1}(\Omega_i))$. Then by Proposition A.8 we obtain the result. ■

Remark A.12. Let $\varepsilon > 0$. With the notations of the above corollary, observe that the map

$$\begin{aligned} L_{\varepsilon, i} : L^2(\Omega) &\rightarrow \mathcal{C}^k([\varepsilon, T], H^{m+2}(\Omega_i)), \\ u_0 &\mapsto (t \mapsto u|_{\Omega_i}(t)), \end{aligned}$$

is continuous for $i = 0, 1$, since [5, Theorem VII.7]

$$\begin{aligned} |u(t)|_{L^2(\Omega)} &\leq |u_0|_{L^2(\Omega)}, \\ |\partial_t u(t)|_{L^2(\Omega)} &= |\nabla \cdot (c \nabla u(t))|_{L^2(\Omega)} \leq \left| \frac{1}{t} \right| |u_0|_{L^2(\Omega)}. \end{aligned}$$

Note that the operator norm of $L_{\varepsilon, i}$ can be bounded uniformly w.r.t. c for $c_{min} \leq c \leq c_{max}$.

Acknowledgement: The authors wish to thank M. Cristofol, Y. Dermenjian, and O. Poisson for numerous discussions on the subject of this article. The authors also thank F. Boyer for helpful comments. The authors also wish to thank an anonymous reviewer for his extremely detailed comments and corrections, which improved the exposition of the article.

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